

$B \rightarrow \pi K, \pi K^*$  and  $\rho K$  Decays :

## CP Violation and Implication for New Physics

Based on: JHEP 0809 (2008) 038.

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## Outline

- Introduction
- Recalculation with dynamical gluon propagator
- The effects of new (pseudo-)scalar couplings
- Summary



## Puzzles in charmless B decays

- Polarization Puzzles
- $\pi\pi$ ,  $\pi K$  Puzzles
- New phase in  $B_s - \bar{B}_s$  mixing
- FB Asymmetry in  $B \rightarrow K^* \mu\mu$

.....

JHEP 0706: 038, 2007

JHEP 0809: 038, 2008

JHEP 0905: 056, 2009

arXiv: 1003. 6051 [hep-ph]

JHEP 1002: 082, 2010

arXiv: 1002. 2758 [hep-ph] (accepted by JHEP)

$B \rightarrow \pi K, \pi K^*$  and  $\rho K$  Decays :

**CP Violation and Implication for New Physics**

**JHEP 0809 (2008) 038**

**Qin Chang, Xin-Qiang Li, Ya-Dong Yang**

These decays are correlated:  
increasing tension  
between Th. & Exp.

### The measurements:

Belle Nature 452 (2008) 332

$$A_{\text{CP}}(B^- \rightarrow K^- \pi^0) \equiv \frac{\Gamma(B^- \rightarrow K^- \pi^0) - \Gamma(B^+ \rightarrow K^+ \pi^0)}{\Gamma(B^- \rightarrow K^- \pi^0) + \Gamma(B^+ \rightarrow K^+ \pi^0)} = +0.07 \pm 0.03 \pm 0.01,$$

$$A_{\text{CP}}(\bar{B}^0 \rightarrow K^- \pi^+) \equiv \frac{\Gamma(\bar{B}^0 \rightarrow K^- \pi^+) - \Gamma(B^0 \rightarrow K^+ \pi^-)}{\Gamma(\bar{B}^0 \rightarrow K^- \pi^+) + \Gamma(B^0 \rightarrow K^+ \pi^-)} = -0.094 \pm 0.018 \pm 0.008.$$

### The difference:

$$\Delta A \equiv A_{\text{CP}}(B^- \rightarrow K^- \pi^0) - A_{\text{CP}}(\bar{B}^0 \rightarrow K^- \pi^+) = \boxed{0.164 \pm 0.037} \quad \text{at } 4.4 \sigma$$

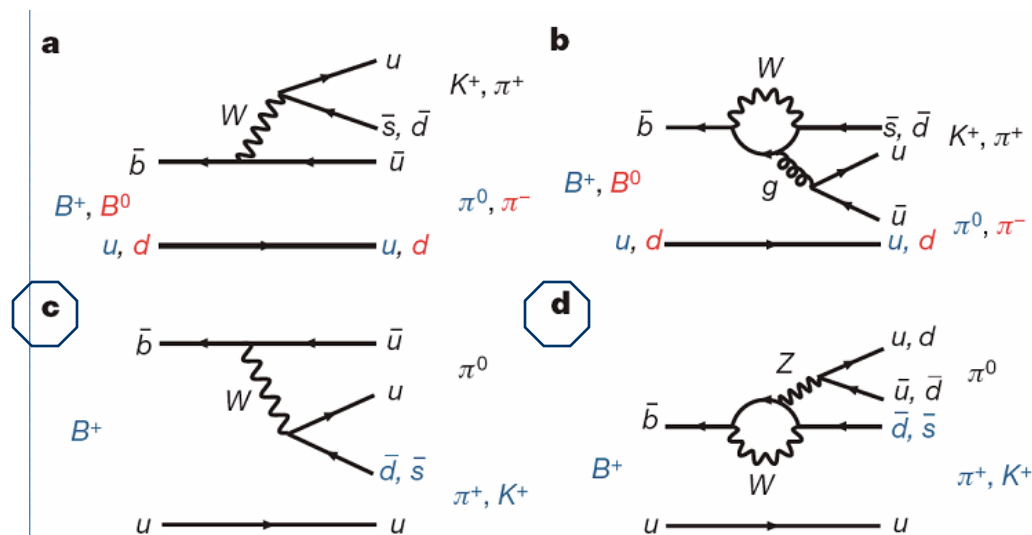
### The average of BABAR, Belle, CDF & CLEO

$$A_{\text{CP}}(B^- \rightarrow K^- \pi^0) = 0.050 \pm 0.025,$$

$$A_{\text{CP}}(\bar{B}^0 \rightarrow K^- \pi^+) = -0.097 \pm 0.012,$$

$$\Delta A = \boxed{0.147 \pm 0.028} \quad \text{at } 5.3 \sigma$$

### The SM expectations:



$$A_{CP}(B^- \rightarrow K^- \pi^0)$$

$$A_{CP}(\bar{B}^0 \rightarrow K^- \pi^+)$$

**Should be very close**

### Theoretical predictions so far

$$\begin{cases} A_{CP}(B_u^- \rightarrow \pi^0 K^-)_{QCDF} = -3.6\%, \\ A_{CP}(\bar{B}_d^0 \rightarrow \pi^+ K^-)_{QCDF} = -4.1\%; \end{cases} \quad QCDF \text{ Scenario S4}$$

$$\begin{cases} A_{CP}(B_u^- \rightarrow \pi^0 K^-)_{PQCD} = (-1_{-5}^{+3})\%, \\ A_{CP}(\bar{B}_d^0 \rightarrow \pi^+ K^-)_{PQCD} = (-9_{-8}^{+6})\%; \end{cases} \quad pQCD$$

**No difference!**

$$\begin{cases} A_{CP}(B_u^- \rightarrow \pi^0 K^-)_{SCET} = (-11 \pm 9 \pm 11 \pm 2)\%, \\ A_{CP}(\bar{B}_d^0 \rightarrow \pi^+ K^-)_{SCET} = (-6 \pm 5 \pm 6 \pm 2)\%. \end{cases} \quad SCET$$

Nucl. Phys. B 675 (2003) 333

Phys. Rev. D 72 (2005) 114005

Phys. Rev. D 74 (2006) 014003

## Possible Implications

The mismatch may be due to:

- ★ Our current limited understanding of strong dynamics involved in hadronic B decays, say, strong phase
- ★ Equally, new physics

### Recalculation with dynamical gluon propagator

#### Scheme I:

$$\int_0^1 \frac{dx}{x} \rightarrow X_A = (1 + \rho_A e^{i\phi_A}) \ln \frac{m_B}{\Lambda_h}, \quad \int_0^1 dy \frac{\ln y}{y} \rightarrow -\frac{1}{2} (X_A)^2$$

$$\rho_A \leq 1 \text{ and } \phi_A \text{ unrestricted}$$

#### Scheme II:

We modify QCDf with dynamical gluon propagator

supported by recent studies with Lattice QCD simulations, Schwinger-Dyson eq.

#### J. M. Cornwall prescription:

$$D(q^2) = \frac{1}{q^2 - M_g^2(q^2) + i\epsilon}$$

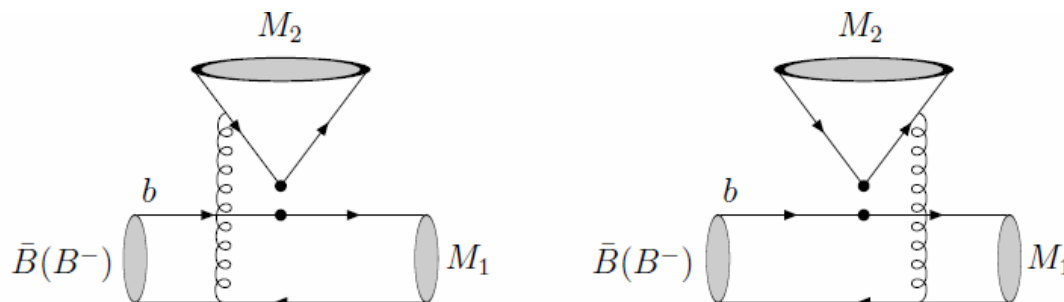
$$\alpha_s(q^2) = \frac{4\pi}{\beta_0 \ln \left( \frac{q^2 + 4M_g^2(q^2)}{\Lambda_{\text{QCD}}^2} \right)}$$

$$M_g^2(q^2) = m_g^2 \left[ \frac{\ln \left( \frac{q^2 + 4m_g^2}{\Lambda_{\text{QCD}}^2} \right)}{\ln \left( \frac{4m_g^2}{\Lambda_{\text{QCD}}^2} \right)} \right]^{-\frac{12}{11}}$$

Phys. Rev. D 26 (1982) 1453;  
Phys. Rev. D 44 (1991) 1285;



### Hard spectator scattering contributions



Space-like

$$H_i(M_1 M_2) = \frac{B_{M_1 M_2}}{A_{M_1 M_2}} \int_0^1 dx dy d\xi \frac{\alpha_s(q^2)}{\xi} \Phi_{B1}(\xi) \Phi_{M_2}(x) \times \left[ \frac{\Phi_{M_1}(y)}{\bar{x}(\bar{y} + \omega^2(q^2)/\xi)} + r_\chi^{M_1} \frac{\phi_{m_1}(y)}{x(\bar{y} + \omega^2(q^2)/\xi)} \right]$$

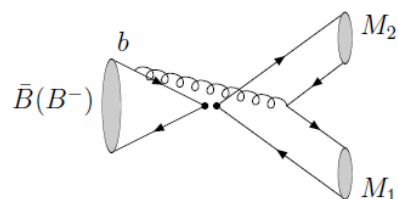
Real

for the contributions of operators  $Q_{i=1-4,9,10}$ ,

$$H_i(M_1 M_2) = -\frac{B_{M_1 M_2}}{A_{M_1 M_2}} \int_0^1 dx dy d\xi \frac{\alpha_s(q^2)}{\xi} \Phi_{B1}(\xi) \Phi_{M_2}(x) \times \left[ \frac{\Phi_{M_1}(y)}{x(\bar{y} + \omega^2(q^2)/\xi)} + r_\chi^{M_1} \frac{\phi_{m_1}(y)}{\bar{x}(\bar{y} + \omega^2(q^2)/\xi)} \right],$$

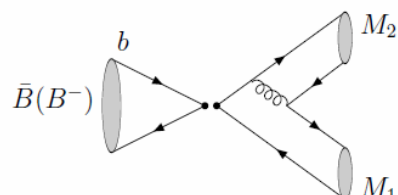
for  $Q_{i=5,7}$ , and  $H_i(M_1 M_2) = 0$  for  $Q_{i=6,8}$ .

### Annihilation contributions

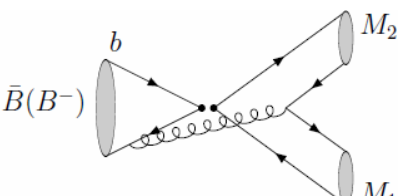


$$A_1^i = \pi \int_0^1 dx dy \alpha_s(q^2) \left\{ \left[ \frac{\bar{x}}{(\bar{x}y - \omega^2(q^2) + i\epsilon)(1 - x\bar{y})} + \frac{1}{(\bar{x}y - \omega^2(q^2) + i\epsilon)\bar{x}} \right] \Phi_{M_1}(y)\Phi_{M_2}(x) + \frac{2}{\bar{x}y - \omega^2(q^2) + i\epsilon} r_\chi^{M_1} r_\chi^{M_2} \phi_{m_1}(y)\phi_{m_2}(x) \right\},$$

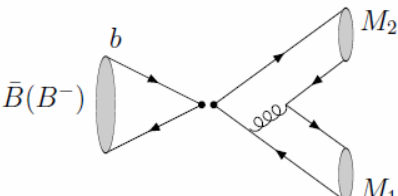
$$A_1^f = A_2^f = 0,$$



$$A_2^i = \pi \int_0^1 dx dy \alpha_s(q^2) \left\{ \left[ \frac{y}{(\bar{x}y - \omega^2(q^2) + i\epsilon)(1 - x\bar{y})} + \frac{1}{(\bar{x}y - \omega^2(q^2) + i\epsilon)y} \right] \Phi_{M_1}(y)\Phi_{M_2}(x) + \frac{2}{\bar{x}y - \omega^2(q^2) + i\epsilon} r_\chi^{M_1} r_\chi^{M_2} \phi_{m_1}(y)\phi_{m_2}(x) \right\},$$



$$A_3^i = \pi \int_0^1 dx dy \alpha_s(q^2) \left\{ \frac{2\bar{y}}{(\bar{x}y - \omega^2(q^2) + i\epsilon)(1 - x\bar{y})} r_\chi^{M_1} \phi_{m_1}(y)\Phi_{M_2}(x) - \frac{2x}{(\bar{x}y - \omega^2(q^2) + i\epsilon)(1 - x\bar{y})} r_\chi^{M_2} \phi_{m_2}(x)\Phi_{M_1}(y) \right\},$$

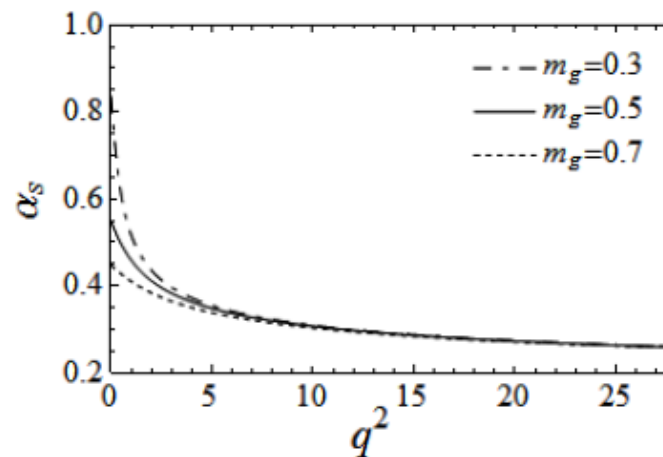
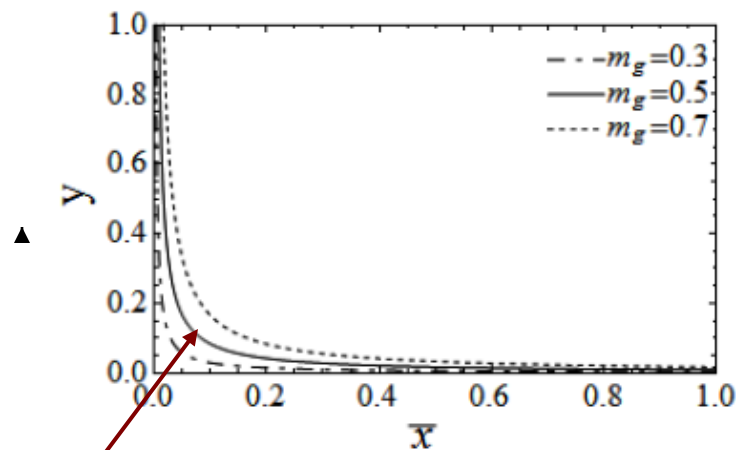


$$A_3^f = \pi \int_0^1 dx dy \alpha_s(q^2) \left\{ \frac{2(1 + \bar{x})}{(\bar{x}y - \omega^2(q^2) + i\epsilon)\bar{x}} r_\chi^{M_1} \phi_{m_1}(y)\Phi_{M_2}(x) + \frac{2(1 + y)}{(\bar{x}y - \omega^2(q^2) + i\epsilon)y} r_\chi^{M_2} \phi_{m_2}(x)\Phi_{M_1}(y) \right\},$$

**Time-like**

**Imaginary**

Typically,  $m_g = 500 \pm 200$  MeV



cancellations

For space-like, it furnishes a cut-off  
 For time-like, it crosses real states,  
 therefore, strong phases

### Results for Branching ratios $\times 10^{-6}$

Decay Mode	QCDF			Experiment data
	$m_g = 0.3$	$m_g = 0.7$	$m_g = 0.45 \sim 0.55$	
$B_u^- \rightarrow \pi^- \bar{K}^0$	44.4	16.8	$23.17 \pm 3.28$	$23.1 \pm 1.0$
$B_u^- \rightarrow \pi^0 K^-$	23.4	9.3	$12.50 \pm 1.65$	$12.9 \pm 0.6$
$\bar{B}_d^0 \rightarrow \pi^+ K^-$	44.7	16.3	$22.71 \pm 3.27$	$19.4 \pm 0.6$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^0$	21.2	7.3	$10.50 \pm 1.63$	$9.9 \pm 0.6$
$B_u^- \rightarrow \pi^- \bar{K}^{*0}$	28.3	5.2	$8.90 \pm 1.59$	$10.0 \pm 0.8$
$B_u^- \rightarrow \pi^0 K^{*-}$	15.2	3.4	$5.25 \pm 0.83$	$6.9 \pm 2.3$
$\bar{B}_d^0 \rightarrow \pi^+ K^{*-}$	28.7	5.3	$9.13 \pm 1.68$	$10.6 \pm 0.9$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^{*0}$	13.4	1.9	$3.89 \pm 0.82$	$2.4 \pm 0.7$
$B_u^- \rightarrow \rho^- \bar{K}^0$	31.8	5.6	$10.27 \pm 1.96$	$8.0^{+1.5}_{-1.4}$
$B_u^- \rightarrow \rho^0 K^-$	14.9	2.5	$4.81 \pm 0.94$	$3.81^{+0.48}_{-0.46}$
$\bar{B}_d^0 \rightarrow \rho^+ K^-$	38.6	8.0	$13.42 \pm 2.31$	$8.6^{+0.9}_{-1.1}$
$\bar{B}_d^0 \rightarrow \rho^0 \bar{K}^0$	21.0	4.8	$7.53 \pm 1.25$	$5.4^{+0.9}_{-1.0}$

$$m_g = \boxed{500 \pm 200 \text{ MeV}} \longleftrightarrow \boxed{500 \pm 50 \text{ MeV}}$$

## The known ratios

Our results

$$R_c \equiv 2 \left[ \frac{Br(B^- \rightarrow \pi^0 K^-)}{Br(B^- \rightarrow \pi^- K^0)} \right] = 1.08 \pm 0.30,$$
$$R_n \equiv \frac{1}{2} \left[ \frac{Br(\bar{B}^0 \rightarrow \pi^+ K^-)}{Br(\bar{B}^0 \rightarrow \pi^0 K^0)} \right] = 1.08 \pm 0.32,$$

Measurements

$$1.12 \pm 0.10$$

$$0.98 \pm 0.09$$

So far, so good.

However, CPA.....

### Results for CPAs $\times 10^{(-2)}$

Decay Mode	QCDF			Experiment data
	$m_g = 0.3$	$m_g = 0.7$	$m_g = 0.45 \sim 0.55$	
$B_u^- \rightarrow \pi^- \bar{K}^0$	0.06	0.19	$0.10 \pm 0.08$	$0.9 \pm 2.5$
$B_u^- \rightarrow \pi^0 K^-$	-11.6	-8.3	$-10.85 \pm 0.84$	$5.0 \pm 2.5$
$\bar{B}_d^0 \rightarrow \pi^+ K^-$	-11.0	-11.4	$-12.38 \pm 0.69$	$-9.7 \pm 1.2$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^0$	2.5	0.1	$1.39 \pm 0.35$	$-14 \pm 11$
$B_u^- \rightarrow \pi^- \bar{K}^{*0}$	0.3	-0.0	$0.16 \pm 0.16$	$-11.4 \pm 6.1$
$B_u^- \rightarrow \pi^0 K^{*-}$	-27.0	-34.1	$-41.20 \pm 6.69$	$4 \pm 29$
$\bar{B}_d^0 \rightarrow \pi^+ K^{*-}$	-27.2	-47.6	$-47.58 \pm 8.42$	$-10 \pm 11$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^{*0}$	3.9	2.1	$4.67 \pm 1.14$	$-9_{-23}^{+32}$
$B_u^- \rightarrow \rho^- \bar{K}^0$	0.1	1.2	$0.53 \pm 0.21$	$-12 \pm 17$
$B_u^- \rightarrow \rho^0 K^-$	28.1	49.7	$46.27 \pm 5.94$	$37 \pm 11$
$\bar{B}_d^0 \rightarrow \rho^+ K^-$	19.3	31.5	$31.40 \pm 4.63$	$15 \pm 13$
$\bar{B}_d^0 \rightarrow \rho^0 \bar{K}^0$	-4.2	0.2	$-3.26 \pm 1.29$	$-2 \pm 29$

 Fine

**In sharp contrast to the Exp., i.e., we can not resolve the difference of CPA**

So far, what we get:

★ nice match for 12 branching ratios with ONE parameter

★ bad match for one CPV

“ $\pi K$  puzzle” still exists

Suspect:

New Physics ?

## • The effects of (pseudo-)scalar couplings

We explore new physics effects in a model independent way

The general four-quark **tensor operators**

$$O_T^q = \bar{s}\sigma_{\mu\nu}(1 + \gamma_5)b \otimes \bar{q}\sigma^{\mu\nu}(1 + \gamma_5)q, \quad O_T^{\prime q} = \bar{s}_i\sigma_{\mu\nu}(1 + \gamma_5)b_j \otimes \bar{q}_j\sigma^{\mu\nu}(1 + \gamma_5)q_i$$



**(pseudo-)scalar operators**

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = \frac{G_F}{\sqrt{2}} \sum_{q=u,d} |V_{tb}V_{ts}^*| e^{i\delta_S^q} \left[ C_{S1}^q O_{S1}^q + C_{S8}^q O_{S8}^q \right] + \text{h.c.}$$

$$O_{S1}^u = \bar{s}(1 + \gamma_5)b \otimes \bar{u}(1 + \gamma_5)u$$

$$O_{S8}^u = \bar{s}_i(1 + \gamma_5)b_j \otimes \bar{u}_j(1 + \gamma_5)u_i$$

$$O_{S1}^d = \bar{s}(1 + \gamma_5)b \otimes \bar{d}(1 + \gamma_5)d$$

$$O_{S8}^d = \bar{s}_i(1 + \gamma_5)b_j \otimes \bar{d}_j(1 + \gamma_5)d_i,$$

**i.e., scalar FCNCs**

In the scanning, all theoretical inputs and Exp. uncertainties are included



## Constraints & Resolutions

We subdivided this NP in five cases:

- Case I:  $b \rightarrow su\bar{u}$  operators  $O_{S1}^u$  and  $O_{S8}^u$ ,
- Case II:  $b \rightarrow sdd\bar{d}$  operators  $O_{S1}^d$  and  $O_{S8}^d$ ,
- Case III:  $b \rightarrow sdd\bar{d}$  operator  $O_{S1}^d$  solely,
- Case IV: only color singlet operators  $O_{S1}^u$  and  $O_{S1}^d$ ,
- Case V: all the operators  $O_{S1}^u$ ,  $O_{S8}^u$ ,  $O_{S1}^d$  and  $O_{S8}^d$ .

### The fitted branching ratios:

Decay Mode	Experiment	NP				
	data	Case I	Case II	Case III	Case IV	Case V
$B_u^- \rightarrow \pi^- \bar{K}^0$	$23.1 \pm 1.0$	—	$23.0 \pm 1.0$	$22.9 \pm 0.9$	$21.5 \pm 0.3$	$22.4 \pm 0.9$
$B_u^- \rightarrow \pi^0 K^-$	$12.9 \pm 0.6$	$12.1 \pm 0.4$	$12.8 \pm 0.7$	$12.7 \pm 0.6$	$12.1 \pm 0.3$	$12.1 \pm 0.4$
$\bar{B}_d^0 \rightarrow \pi^+ K^-$	$19.4 \pm 0.6$	$20.2 \pm 0.3$	—	—	$20.4 \pm 0.2$	$20.1 \pm 0.4$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^0$	$9.9 \pm 0.6$	$9.0 \pm 0.3$	$9.9 \pm 0.6$	$10.0 \pm 0.7$	$9.0 \pm 0.2$	$9.1 \pm 0.4$
$B_u^- \rightarrow \pi^0 K^{*-}$	$6.9 \pm 2.3$	$4.2 \pm 0.2$	$4.4 \pm 0.4$	$4.4 \pm 0.4$	$4.3 \pm 0.3$	$4.3 \pm 0.3$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^{*0}$	$2.4 \pm 0.7$	$3.4 \pm 0.3$	$3.5 \pm 0.2$	$3.5 \pm 0.2$	$3.1 \pm 0.3$	$2.9 \pm 0.2$
$B_u^- \rightarrow \rho^- \bar{K}^0$	$8.0^{+1.5}_{-1.4}$	—	$8.6 \pm 0.7$	$8.6 \pm 0.7$	$7.4 \pm 0.4$	$7.1 \pm 0.4$
$B_u^- \rightarrow \rho^0 K^-$	$3.81^{+0.48}_{-0.46}$	$3.4 \pm 0.2$	—	—	$3.4 \pm 0.2$	$3.4 \pm 0.2$
$\bar{B}_d^0 \rightarrow \rho^+ K^-$	$8.6^{+0.9}_{-1.1}$	$9.7 \pm 0.5$	—	—	$9.7 \pm 0.5$	$9.8 \pm 0.5$
$\bar{B}_d^0 \rightarrow \rho^0 \bar{K}^0$	$5.4^{+0.9}_{-1.0}$	—	$6.5 \pm 0.4$	$6.5 \pm 0.4$	$5.5 \pm 0.3$	$5.4 \pm 0.4$

To leading order,

$B_u^- \rightarrow \pi^- \bar{K}^{*0}$  and  $\bar{B}_d^0 \rightarrow \pi^+ \bar{K}^{*-}$  decays do not receive these NP contributions

### The direct CP asymmetries:

Only Br as inputs

Decay Mode	Experiment data	NP				
		Case I	Case II	Case III	Case IV	Case V
$B_u^- \rightarrow \pi^- \bar{K}^0$	$0.9 \pm 2.5$	—	$1.7 \pm 2.9$	$2.0 \pm 0.2$	$3.9 \pm 1.0$	$3.2 \pm 1.3$
$B_u^- \rightarrow \pi^0 K^-$	$5.0 \pm 2.5$	$8.8 \pm 6.4$	$1.1 \pm 0.9$	$1.2 \pm 0.9$	$2.8 \pm 5.5$	$1.8 \pm 1.3$
$\bar{B}_d^0 \rightarrow \pi^+ K^-$	$-9.7 \pm 1.2$	$-5.7 \pm 4.4$	—	—	$-10.0 \pm 0.8$	$-9.2 \pm 1.3$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^0$	$-14 \pm 11$	$-18.6 \pm 7.5$	$-12.8 \pm 3.9$	$-12.6 \pm 1.6$	$-10.2 \pm 7.0$	$-8.2 \pm 2.8$
$B_u^- \rightarrow \pi^0 K^{*-}$	$4 \pm 29$	$4.2 \pm 19.3$	$-8.1 \pm 3.3$	$-8.0 \pm 3.3$	$-4.9 \pm 19.7$	$-13.2 \pm 4.6$
$\bar{B}_d^0 \rightarrow \pi^0 \bar{K}^{*0}$	$-9_{-23}^{+32}$	$-61.7 \pm 22.0$	$-49.9 \pm 3.4$	$-49.8 \pm 3.8$	$-52.8 \pm 24.2$	$-47.0 \pm 6.5$
$B_u^- \rightarrow \rho^- \bar{K}^0$	$-12 \pm 17$	—	$-5.9 \pm 10.9$	$-6.5 \pm 0.8$	$-15.1 \pm 4.2$	$-13.1 \pm 5.9$
$B_u^- \rightarrow \rho^0 K^-$	$37 \pm 11$	$32.8 \pm 16.5$	—	—	$48.3 \pm 3.5$	$43.9 \pm 5.2$
$\bar{B}_d^0 \rightarrow \rho^+ K^-$	$15 \pm 13$	$19.2 \pm 12.9$	—	—	$31.9 \pm 2.7$	$28.0 \pm 4.1$
$\bar{B}_d^0 \rightarrow \rho^0 \bar{K}^0$	$-2 \pm 29$	—	$-8.1 \pm 8.1$	$-8.5 \pm 0.9$	$-14.9 \pm 3.0$	$-13.5 \pm 4.4$

Case IV

Theo.  $\Delta A = 0.128 \pm 0.056$

$$\begin{matrix} C_{S1}^u \\ C_{S8}^u \\ \delta_S^u \end{matrix}$$

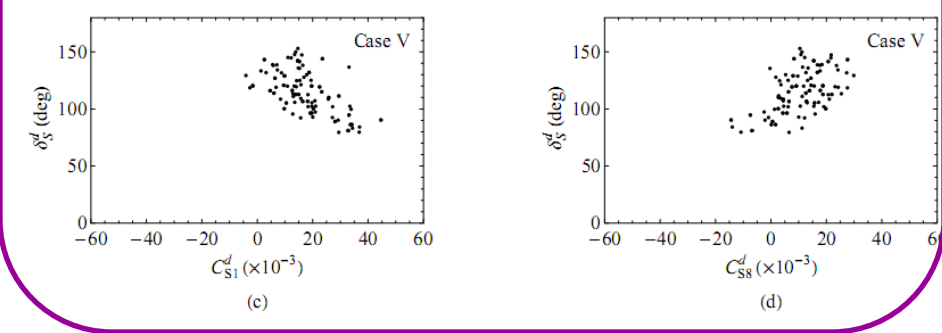
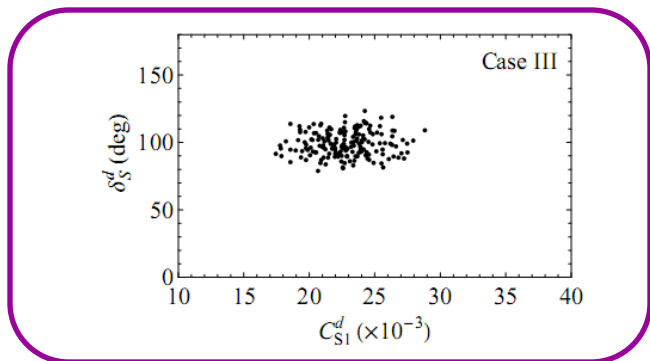
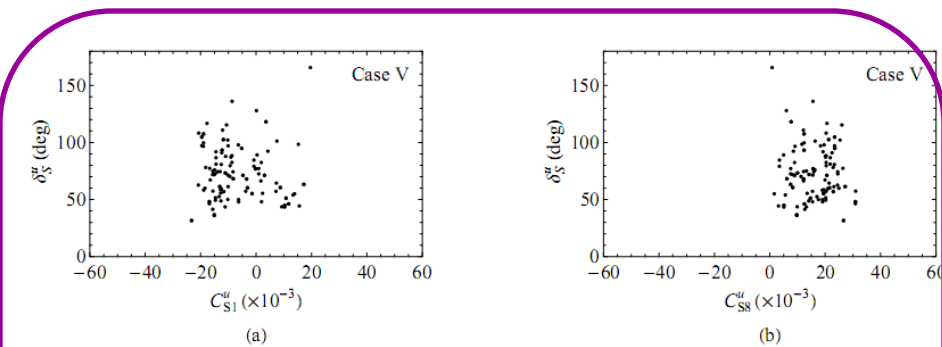
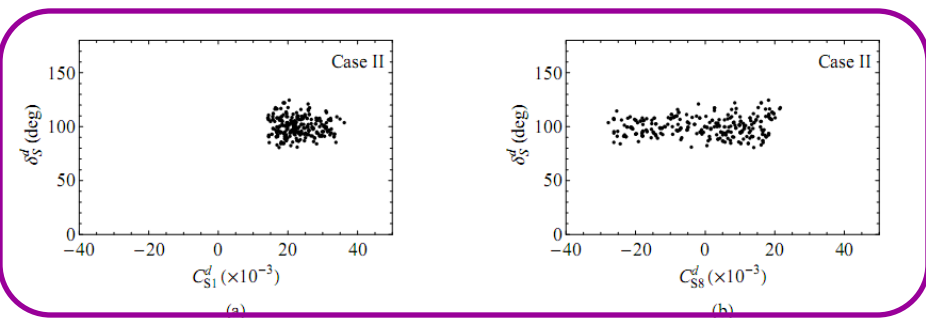
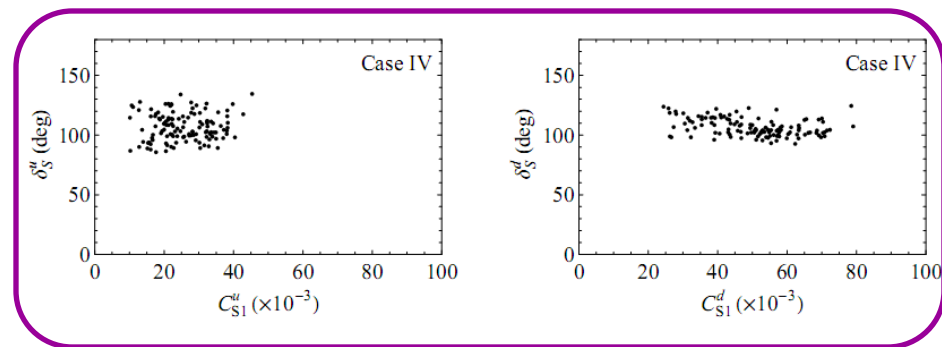
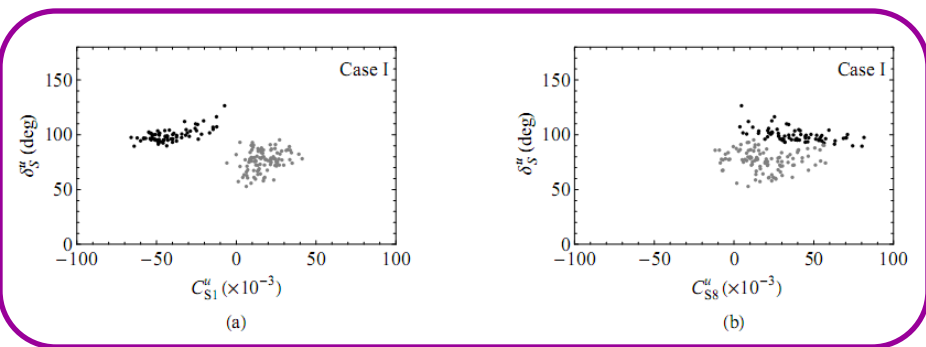
$$\begin{matrix} C_{S1}^d \\ C_{S8}^d \\ \delta_S^d \end{matrix}$$

$$C_{S1}^d$$

$$\begin{matrix} C_{S1}^u \\ C_{S1}^d \end{matrix}$$

Exp.  $\Delta A = 0.147 \pm 0.028$

### NP parameters' space



### Numerical results

NP para.	Case I	Case II	Case III	Case IV	Case V
$C_{S1}^u (\times 10^{-3})$	$-41.6 \pm 13.4$	—	—	$25.8 \pm 8.4$	$-6.7 \pm 10.5$
$C_{S8}^u (\times 10^{-3})$	$38.7 \pm 18.2$	—	—	—	$16.0 \pm 7.1$
$\delta_S^u$	$99.5^\circ \pm 6.1^\circ$	—	—	$107.0^\circ \pm 11.5^\circ$	$73.0^\circ \pm 23.8^\circ$
$C_{S1}^d (\times 10^{-3})$	—	$23.0 \pm 5.1$	$22.8 \pm 2.3$	$50.3 \pm 12.8$	$17.5 \pm 10.1$
$C_{S8}^d (\times 10^{-3})$	—	$-0.8 \pm 13.7$	—	—	$10.5 \pm 9.4$
$\delta_S^d$	—	$100.0^\circ \pm 8.7^\circ$	$99.3^\circ \pm 9.2^\circ$	$106.6^\circ \pm 7.3^\circ$	$114.7^\circ \pm 18.6^\circ$

$$|C_{S1}^u| \approx |C_{S8}^u|$$

**Exotic !**

$$|C_{S1}^d| > |C_{S8}^d|$$

**Color singlet  
dominated !**

$$C_{S1}^d \approx 2 \times C_{S1}^u$$

**Down type  
dominated !**

**Nontrivial NP weak phase!**

## Prediction

The mixing induced CPA in  $\bar{B}^0 \rightarrow \pi^0 K_S$  and  $\rho^0 K_S$

- insensitive to strong phase, more accurate in QCDF
- sensitive to New Physics (new weak phase)
- should be  $\sin(2\beta)_{\Psi K_S}$  in the SM

$$A_f(t) = S_f \sin(\Delta m_d t) - C_f \cos(\Delta m_d t)$$

$$S_f = \mathcal{A}_{\text{CP}}^{\text{mix}}$$

$$-C_f \equiv \mathcal{A}_{\text{CP}}$$

$$\mathcal{A}_{\text{CP}}^{\text{mix}}(\bar{B}^0 \rightarrow f) = \frac{2\text{Im}\lambda_f}{1 + |\lambda_f|^2}$$

$$\lambda_f = -e^{-2i\beta} \bar{A}^{00} / A^{00}$$

$$\sin(2\beta) = \sin(2\beta)_{\Psi K_S} = 0.68 \pm 0.03$$

( $\times 10^{-2}$ )

Decay Mode	Experiment data	SM	NP				
			Case I	Case II	Case III	Case IV	Case V
$\bar{B}_d^0 \rightarrow \pi^0 K_S$	$38 \pm 19$	$77 \pm 4$	$45 \pm 11$	$56 \pm 5$	$57 \pm 3$	$59 \pm 9$	$62 \pm 8$
$\bar{B}_d^0 \rightarrow \rho^0 K_S$	$61^{+25}_{-27}$	$66 \pm 3$	—	$61 \pm 6$	$61 \pm 3$	$56 \pm 3$	$57 \pm 4$

Old data at that time

### ICHEP 08

BABAR  $A_{CP}^{\text{mix}}(\bar{B}^0 \rightarrow \pi^0 K_S) = 0.55 \pm 0.20 \pm 0.03$   
 Belle  $A_{CP}^{\text{mix}}(\bar{B}^0 \rightarrow \pi^0 K_S) = 0.67 \pm 0.31 \pm 0.08$

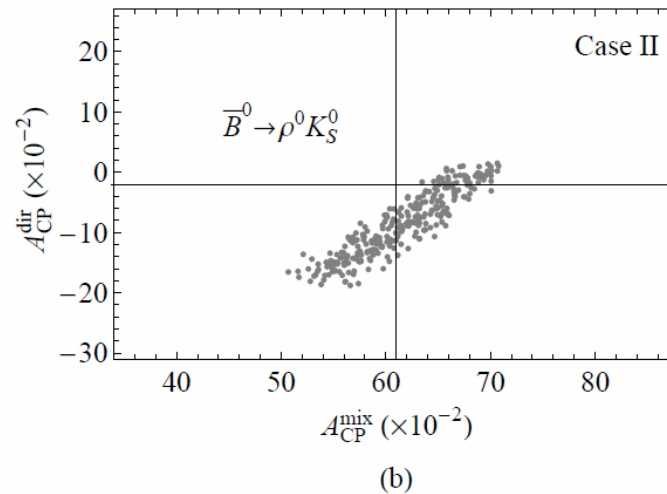
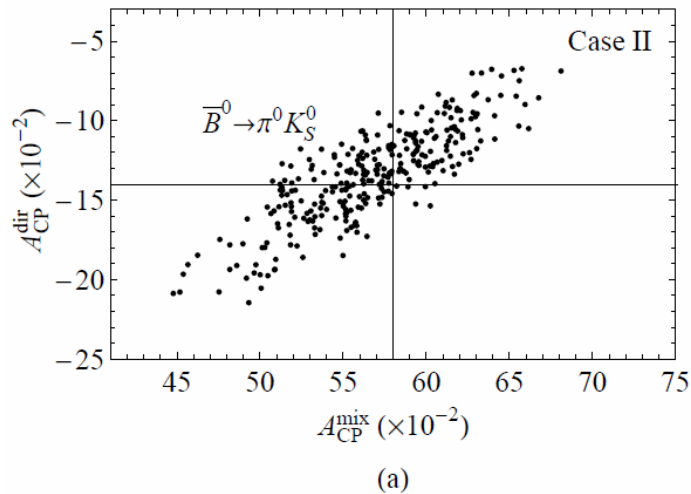
$0.58 \pm 0.17$

### Now, HFAG

$A_{CP}^{\text{mix}}(\pi^0 K_S) = 0.57 \pm 0.17$   
 $A_{CP}^{\text{mix}}(\rho^0 K_S) = 0.54^{+0.18}_{-0.21}$

Compared with the new data, **unexpected match!**

## Correlations:



- The correlations could be important to sign out new physics
- Need more accurate mixing-induced CP to confirm NP suspects



## ▪ Summary

- The knowns of QCD from Lattice and SD should be incorporated into the approaches for B decays
- Dynamical gluon mass could, at least, furnish the natural regulator of end-point div.
- Correlated decays point to NP in the EW penguin sector
- (pseudo-)scalar couplings are helpful for resolving " $\pi K$  puzzle"
- More accurate measurements, especially mixing-induced CPV, are also needed to confirm or refute the NP hints

**Thank you  
for  
your attention!**