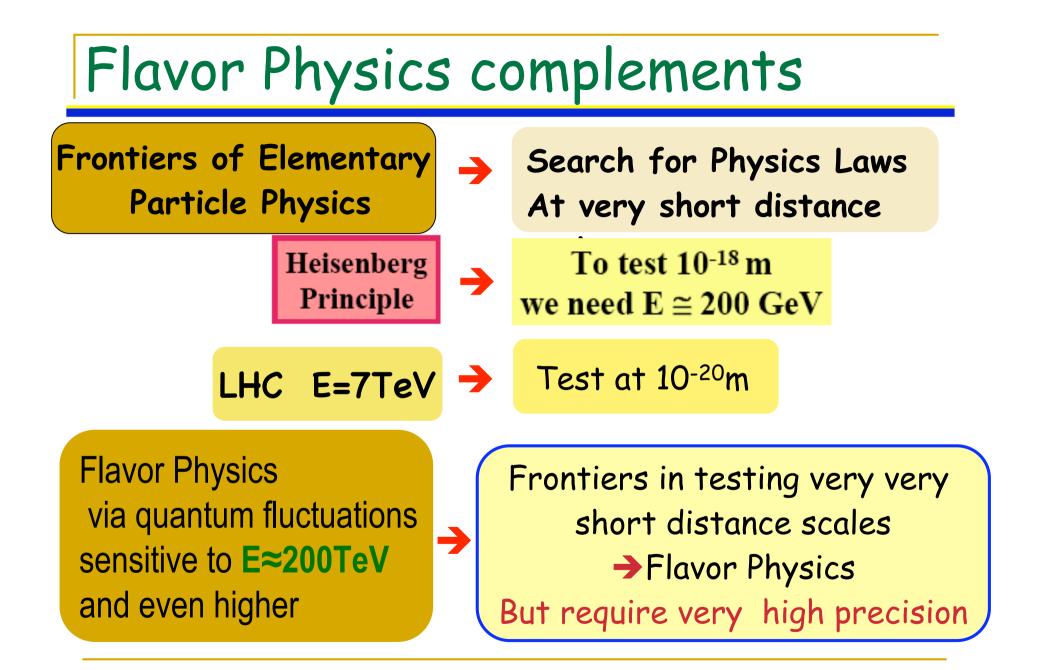
# Recent Progress in Heavy Quark Flavor Physics

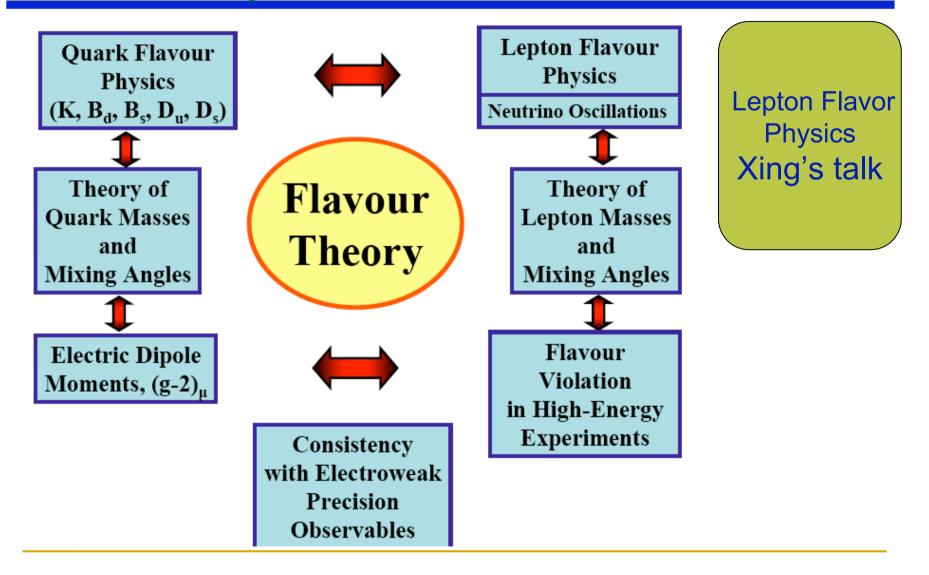
杨亚东 华中师大粒子所

# Outline

- Overview
- Experimental progress
- Puzzling results
- Theoretical progress
- Summary

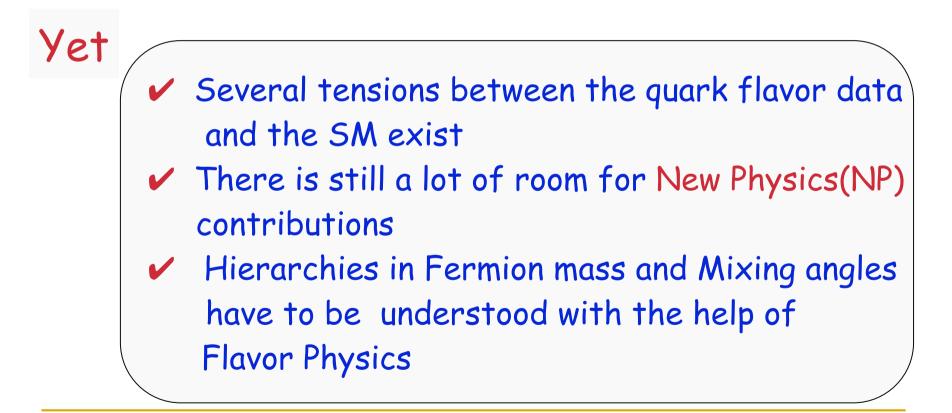


# Flavor Physics-

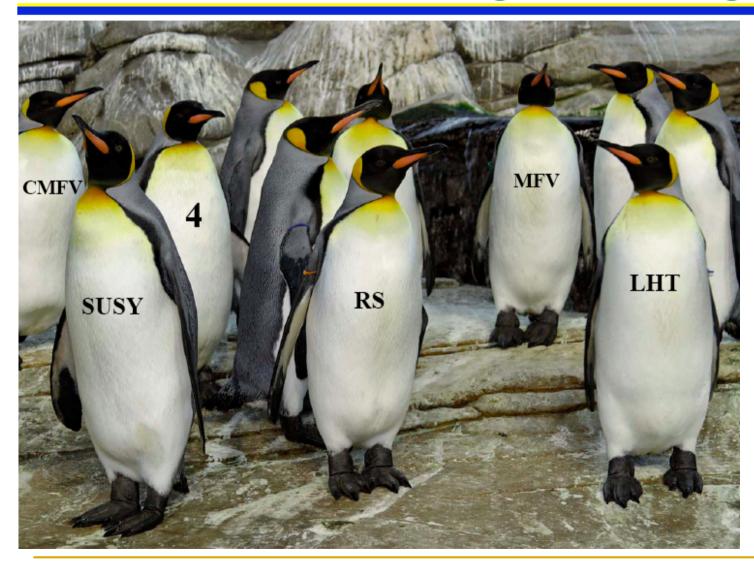


# Quark Flavor Physics

So far, confirmed success of the CKM picture of Flavor Changing Neutral Current of GIM picture

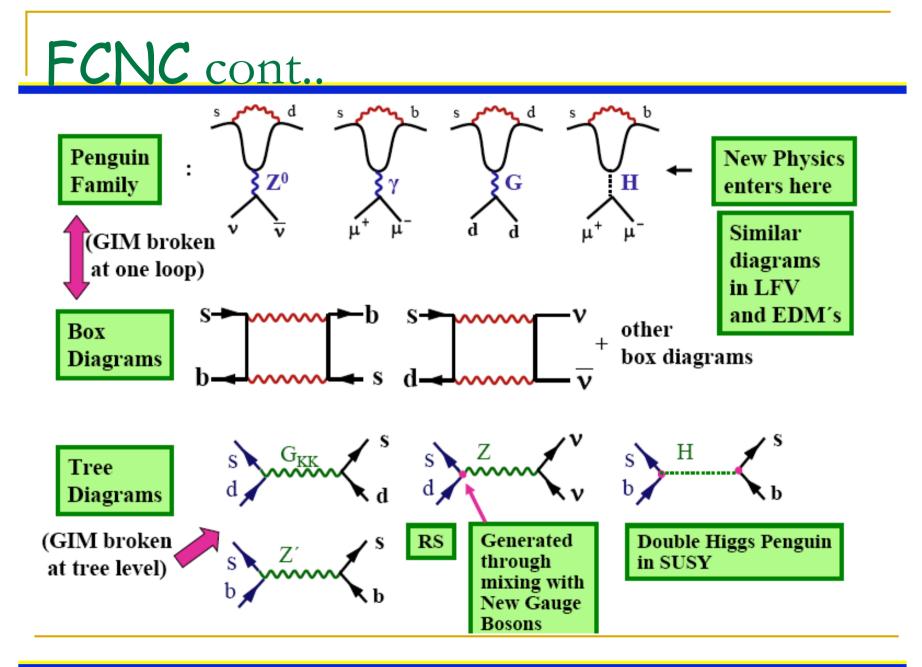


# NP scenarios dancing with Penguins

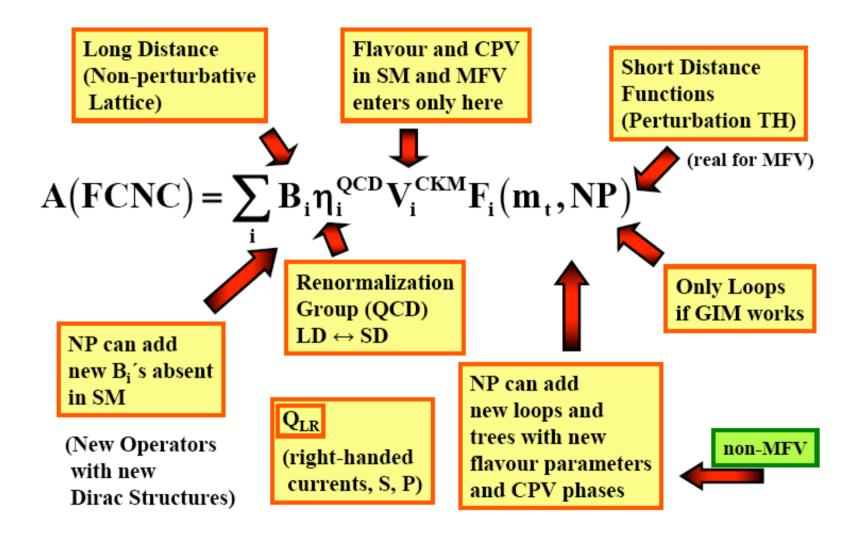


Their flavor structures, parameter spaces, new particles are constrained very much by Penguin processes

#### FCNC



# FCNC cont. Master formula



# FCNC cont.

- Rare but not so rare: precise measurement
   @BaBar, Belle, CDF, D0, LHCb
- Sensitive to New Physics, thus tough constraints on NP scenarios
- Strong correlated & complementary to the direct searches via high energy processes
- Hints for NP have been shown

Groups involved in Heavy Flavor Physics

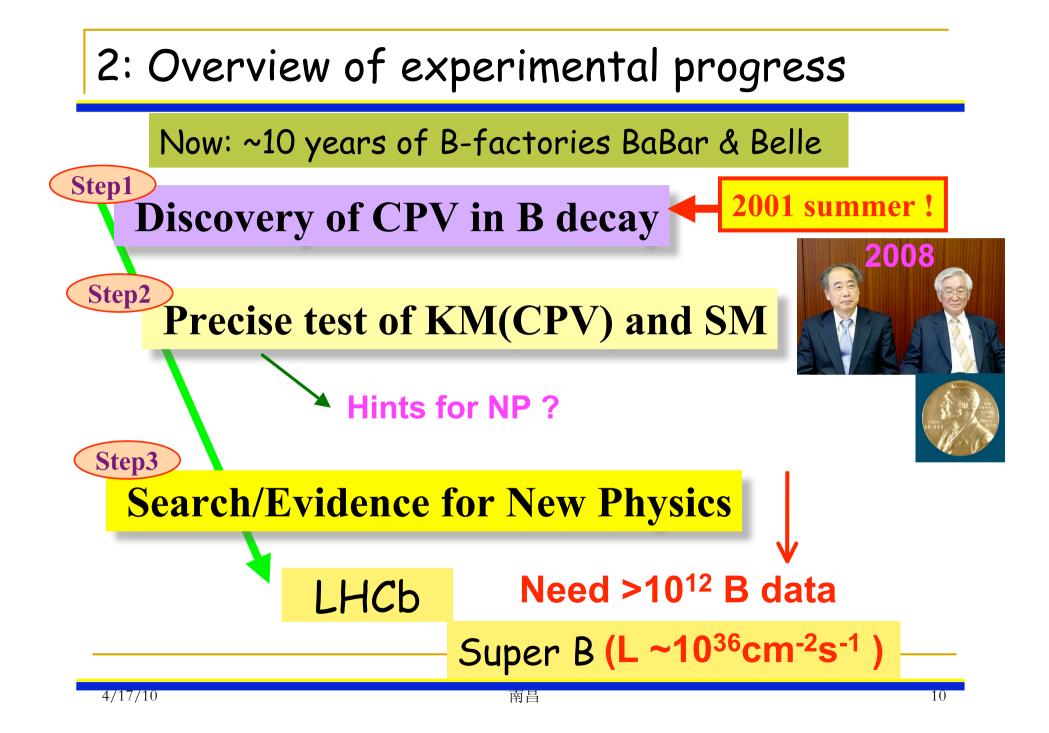
From North China to South China

```
HIT, Dalian(LNNU+DLUT)
Tsinghua, PKU, ITP, IHEP,GUCAS Nankai,
Huabei, Yantai, Henan NU,
USTC
```

Chongqin U., Nanjing Normal, Zhejiang U,

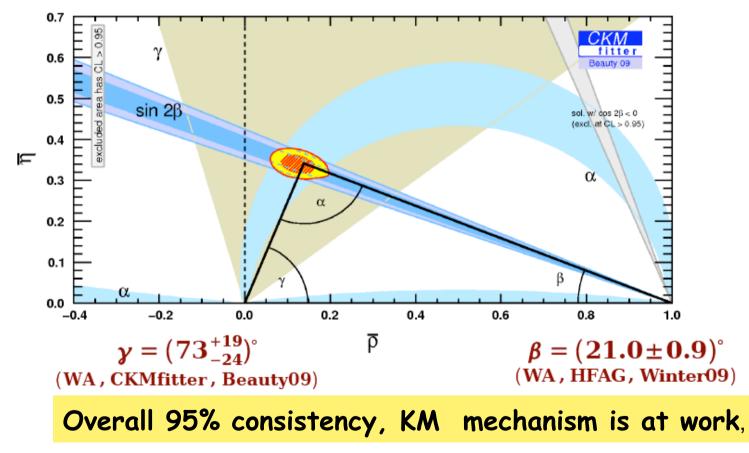
Shanghai, CCNU, Nanchang.....

## So, there is a big asymmetry

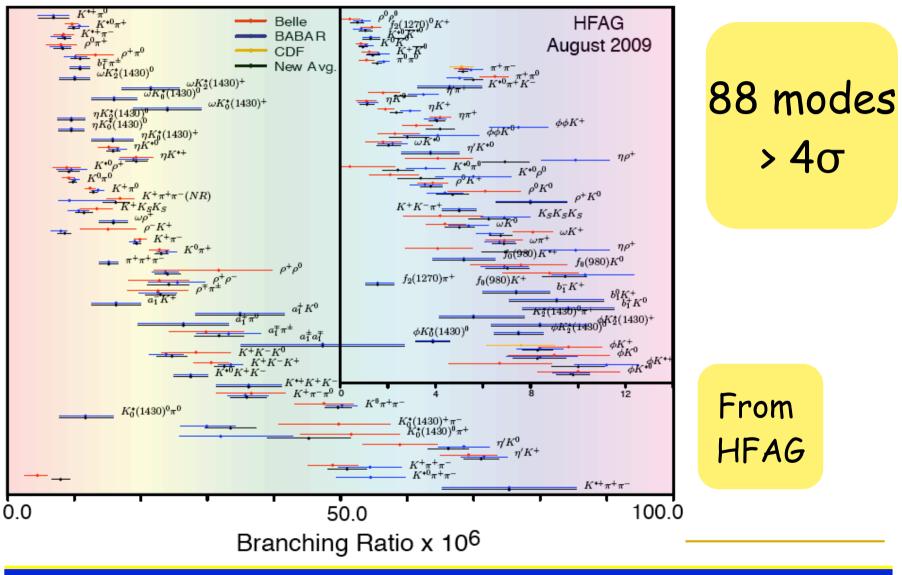


### The unitary triangle

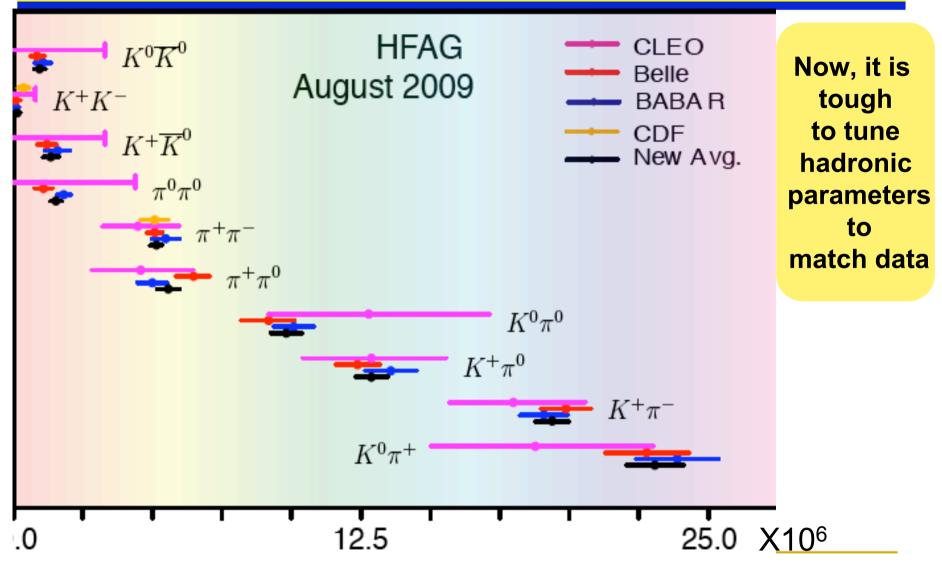
 $\alpha = (89.0^{+4.4}_{-4.2})^{\circ}$ (WA, CKMfitter, Winter09)



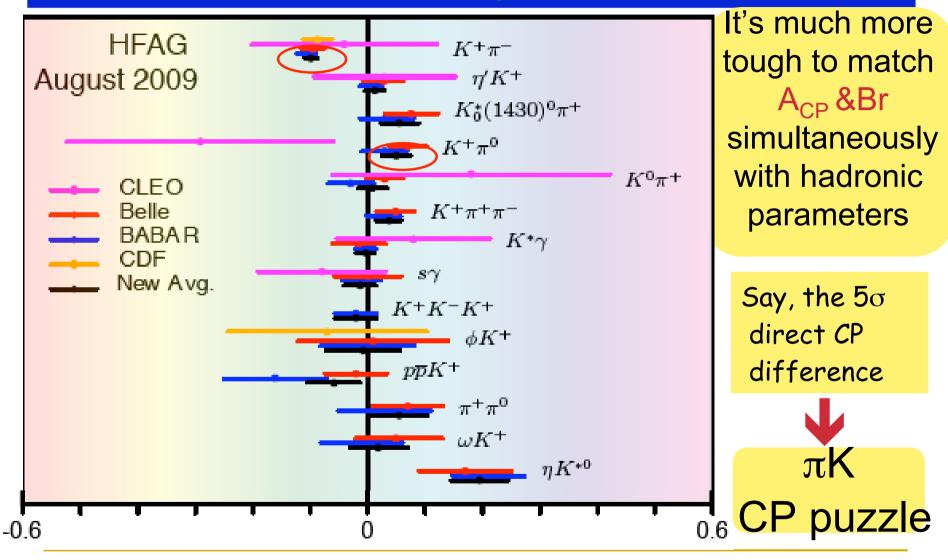
# 2.1 Charmless nonleptonic decays(exp)



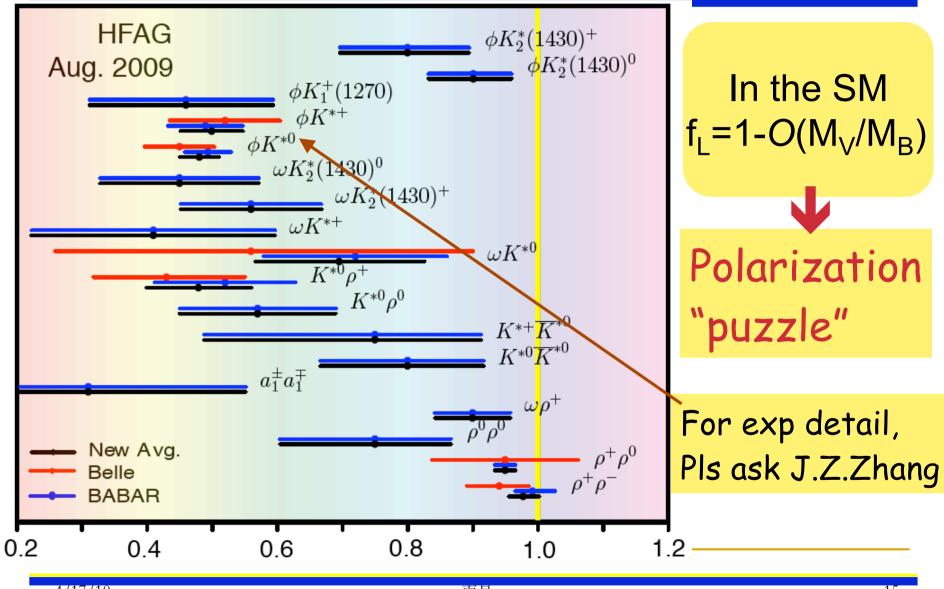
### 2.1 Precise measurements



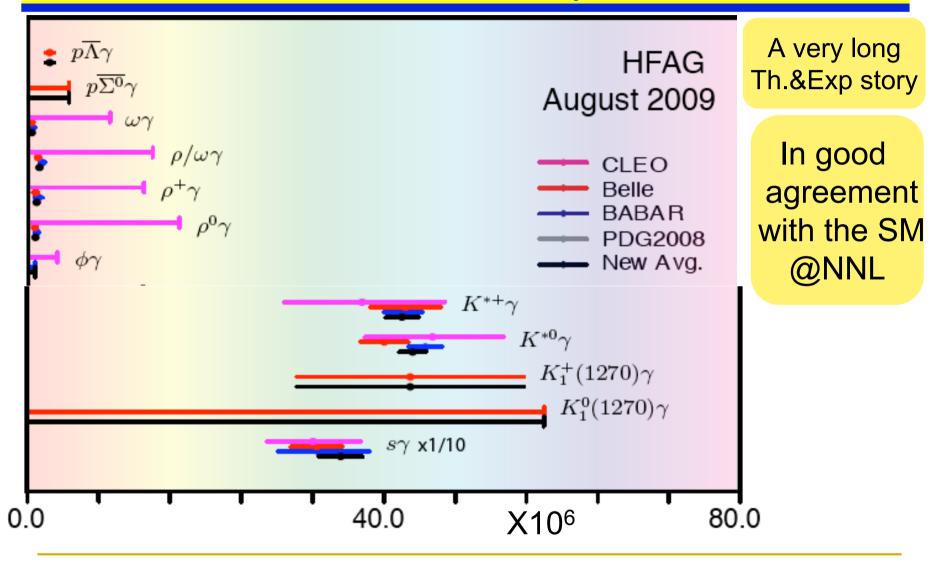
# 2.1 cont. The CP Asymmettries



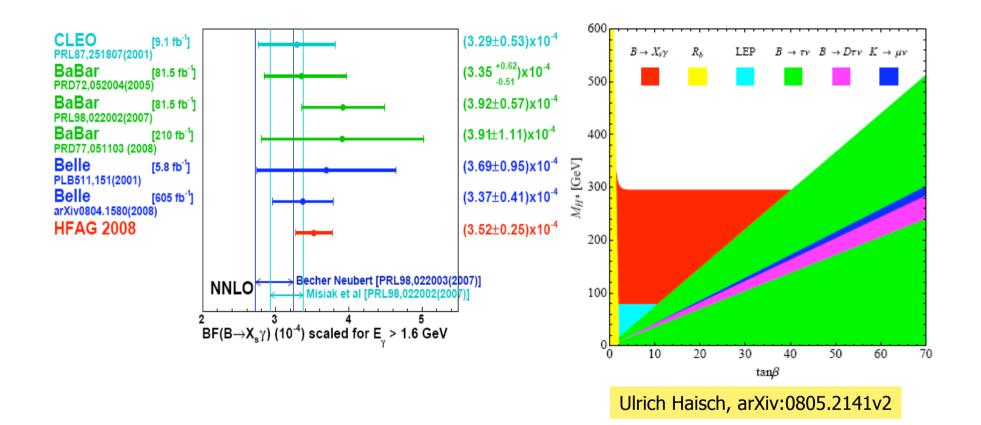
# 2.1 cont. Polarizations in $B \rightarrow VV$



# 2.2 The radiative B decays

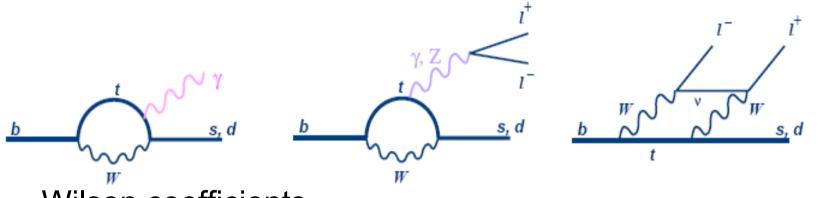


 $B \rightarrow X_{c}\gamma$ 



Together with  $(g-2)_{\mu}$ ,  $B \rightarrow X_s \gamma$  serve as strong constraints on NP

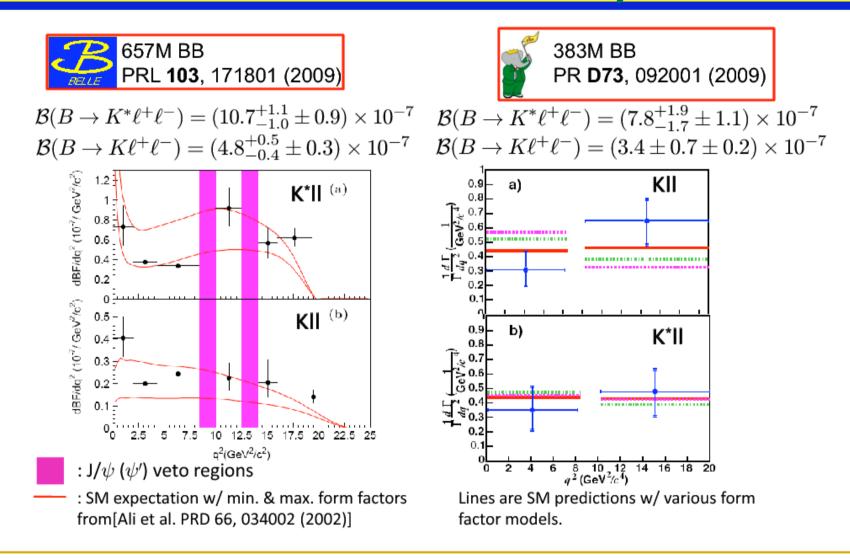
# 2.3 The $A_{FB}$ in $B \rightarrow X_s / / / , K^{(*)} / / / -$

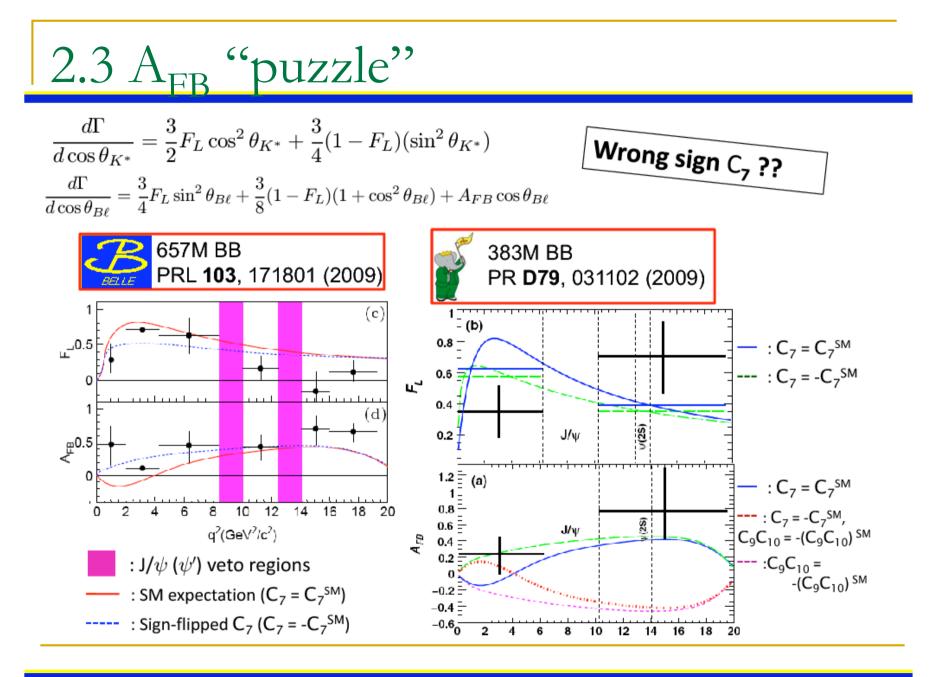


- Wilson coefficients
  - C<sub>7</sub>: from electromagnetic penguin diagram
     C<sub>9</sub>: from vector electroweak
  - $C_{10}$ : from axial vector electroweak
- Differential branching fraction (B.F.) and Forward-backward asymmetry (A<sub>FB</sub>) in B → K<sup>\*</sup>I<sup>+</sup>I<sup>-</sup>

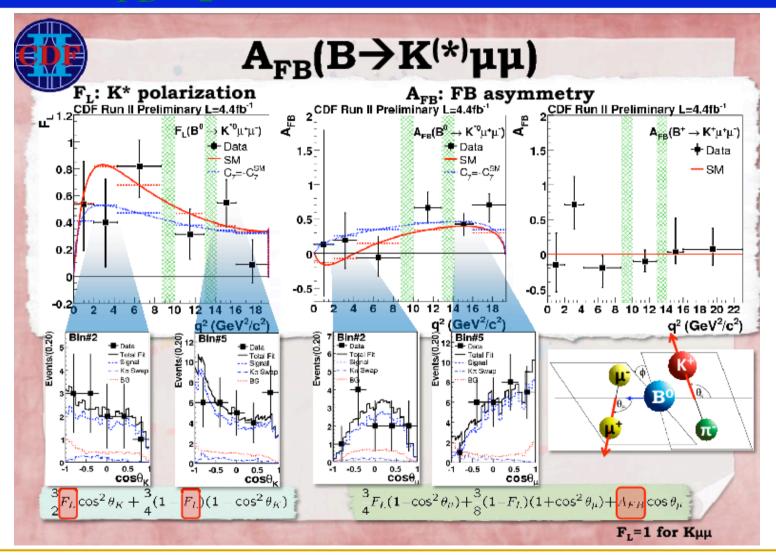
$$\frac{dA_{FB}}{dq^2} \propto -\Re e[\widetilde{C}_9 \widetilde{C}_{10} V A_1 + \frac{M_B m_b}{q^2} \widetilde{C}_7 \widetilde{C}_{10} (V T_2 \cdot (1 - \frac{m_{K^*}}{M_B}) + A_1 T_1 \cdot (1 + \frac{m_{K^*}}{M_B}))]$$

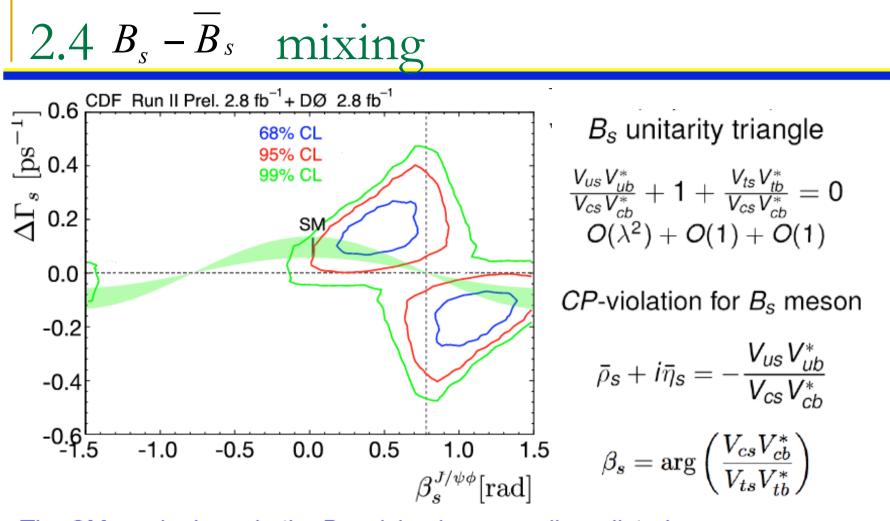
### 2.3 Rare electroweak B decays





# 2.3 A<sub>FB</sub> "puzzle" contd





The SM weak phase in the Bs mixing is very well predicted:

$$\beta_s = (1.035 + 0.049)_{-0.046} \text{ deg} \qquad [\beta_s = 22.3 + 10.2_{-8.0} \text{ deg (Tevatron)}]$$

- There are so many interesting pro
- X, Y, Z states [S.Zhu, Q. Zhao, K.T. Chao, X.Liu, Z.G. Wang...]
- Charmonium production [state of the art calculations by K.T.Chao, J.X.Wang; C.Qiao,Y.Jia, D.Yang......]
- **f**<sub>Ds</sub> **puzzle** [X.Q.Li, Z.Wei, H.b.Li.....]
- I will skip charm physics, since there is a nice review

**Charm Physics:** 

A Field Full with Challenges and Opportunities [X.Q.Li, Z.T.Wei, Front.Phys.China 4:49-74,2009]

# 3. Theories.

- Hadronic dynamics (nonperturbative QCD) complicated and hindered precise Th estimations from QCD very very much.
- We rely on Factorization to separate long and short distance effects
- **LD**: modeled, extracted from exp., Lattice
- SD: well defined, EW phys, QCD loop corrections, RGE.....

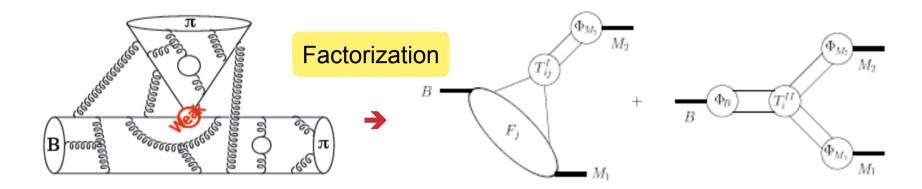
### 3. Charmless B decays: Theory shopping list

- QCD Factorizations [BBNS, PRL'99, NPB'00]
- PQCD [Y.Y.Keum,H.n.Li,A.Sanda, PRD'01, C.Lu, K.Ukai, M.Z.Yang, PRD'01]
- Soft Collinear Effective Theory[SCET]
   [C.W.Bauer et al., PRD'01, '02, '04. J.Chay&C.Kim, PRD'03, NPB'04]
- Transverse momentum dependent [TMD] [J.P.Ma,Q.Wang,PLB613(05)39, JHEP0601, PLB674(09)176]

### Six Quark H<sub>eff</sub> QCDF

[Y.L.Wu, et al., Int.J.Mod.Phys.A25:69-111,2010.]

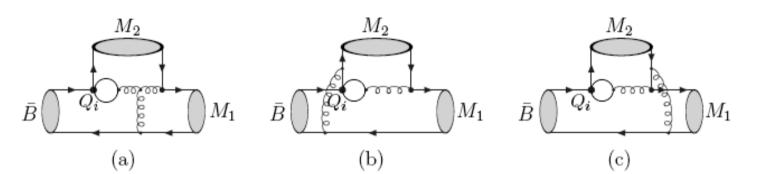
etc

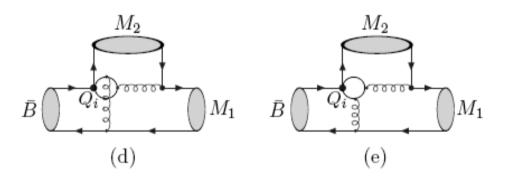


$$\langle M_1 M_2 | C_i O_i | \bar{B} \rangle_{\mathcal{L}_{eff}} = \sum_{\text{terms}} C(\mu_h) \times \left\{ F_{B \to M_1} \times \underbrace{T^{I}(\mu_h, \mu_s)}_{1 + \alpha_s + \dots} \star f_{M_2} \Phi_{M_2}(\mu_s) \right\}$$

$$+f_B\Phi_B(\mu_s)\star\left[\underbrace{T^{\mathrm{II}}(\mu_h,\mu_I)}_{1+\ldots}\star\underbrace{J^{\mathrm{II}}(\mu_I,\mu_s)}_{\alpha_s+\ldots}\right]\star f_{M_1}\Phi_{M_1}(\mu_s)\star f_{M_2}\Phi_{M_2}(\mu_s)\Big\}$$

 $+ 1/m_b$ -suppressed terms





+ bsg ones

• Strong Penguin [partially by X.Q.Li YD, PRD72(05) 074007, 73(06)114027]

#### SCET was embedded in QCDF

#### Spectator -scattering

- One loop J<sup>II</sup>: Becher et al.'04, Beneke, D.S.Yang'05
- One loop T<sup>II</sup> tree amplitudes: Beneke, Jager'05, Pilipp'07
- One loop T<sup>II</sup> penguin amplitudes: Beneke, Jager'06

#### Vertex term

Two loop T<sup>II</sup> tree amplitudes: G.Bell, NPB 822(09)172.
Beneke, Huber,,X.Q.Li,NPB832(10)109
Two loop T<sup>II</sup> penguin amplitudes: Beneke+Li+? In progress

#### Steady progresses!

Approx 70 two-loop vertex diagrams

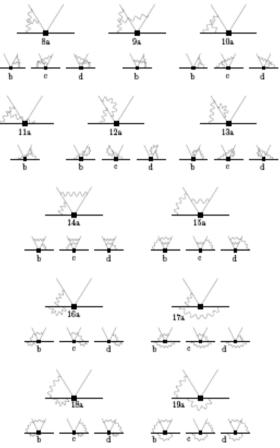
QCD to SCET matching calculation

$$Q_i = \sum_a \tilde{H}_{ia} \tilde{\mathcal{O}}_a$$

with non-local SCET<sub>I</sub> operators  $\tilde{\mathcal{O}}_a$ .

2-loop techniques: Xinqiang's talk@ 18上午

See also Beneke, Huber, Li, NPB811(2009)77



Numerical result (tree amplitudes)

$$a_{1}(\pi\pi) = 1.009 + [0.023 + 0.010i]_{\text{NLO}} + [0.026 + 0.028i]_{\text{NNLO}} - \left[\frac{r_{\text{sp}}}{0.485}\right] \left\{ [0.015]_{\text{LOsp}} + [0.037 + 0.029i]_{\text{NLOsp}} + [0.009]_{\text{tw3}} \right\} = 1.00 + 0.01i \rightarrow 0.93 - 0.02i \quad (\text{if } 2 \times r_{\text{sp}}) \qquad r_{\text{sp}} = \frac{9f_{M_{1}}\hat{f}_{B}}{m_{b}f_{B}^{B\pi}(0)\lambda_{B}}$$

$$a_{2}(\pi\pi) = 0.220 - [0.179 + 0.077i]_{\text{NLO}} - [0.031 + 0.050i]_{\text{NNLO}} \\ + \left[\frac{r_{\text{sp}}}{0.485}\right] \left\{ [0.123]_{\text{LOsp}} + [0.053 + 0.054i]_{\text{NLOsp}} + [0.072]_{\text{tw3}} \right\} \\ = 0.26 - 0.07i \rightarrow 0.51 - 0.02i \quad (\text{if } 2 \times r_{\text{sp}})$$

- The NNLO corrections to the vertex term and spectator scattering are significant individually (about 25% for *a*<sub>2</sub>). But both tend to cancel too bad!
- Largest uncertainty is input parameter dependence: Allows  $|C/T|_{\pi\pi} \approx 0.7$ , if  $\lambda_B$  is small. The colour-suppressed amplitudes are probably dominated by spectator-scattering. But arg  $(C/T_{\pi\pi}) \lesssim 15^{\circ}$ .
- Perturbation theory works at scale  $m_b$  and  $\sqrt{m_b\Lambda}$ . No indication of further large radiative corrections.

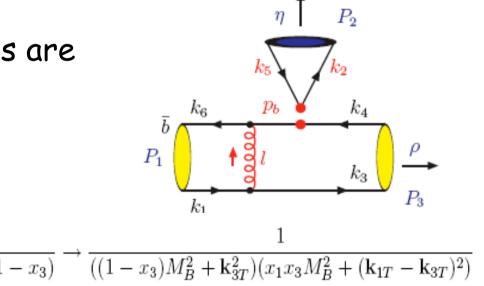
# Theory continued [pQCD]

- $\bullet$  Introduce  $K_{T}$  to kill end-point divergence in spectator-scattering kernel
- Sudakov factor to suppress long distance contributions

Many works by: Caidian, Maozhi, Xiao, Zhentao,

 One loop QCD corrections are in progress

Z.Xiao, PRD08+...



Comparison: QCDF-pQCD-SCETThe end-point issuepQCD:
$$\frac{1}{m_b^2 \bar{x} - k_T^2 + i0}$$
singularity regulated by  $k_T$ 

**BBNS:** Introduce hadronic parameters 
$$\int_0^1 dx/x \to X_A$$
  
 $X_A = (1 + \rho_A e^{i\phi_A}) \ln(m_B/500 \,\text{MeV})$ 

#### SCET:

The annihilation singularity has to do with a potential double counting Arnesen et.al. Same QCD topology appears twice.

# Comparison: QCDF-pQCD-SCET

	SCET	QCDF	pQCD	
Expansion in α <sub>s</sub> (μ <sub>i</sub> )?	No	Yes	Yes	
T, P if Singular convolution	N/A	New parameters	uses k <sub>T</sub>	
Annihilation	Real at "LO", complex "NLO"	Complex, new parameters	perturbative, large phases	
Charm Loop?	Non- perturbative	Perturbative		
Number of fit parameters	Most	Middle	N/A	

Symmetry approaches

To reduce parameters, one has

- SU(2) isospin symmetry for pions
- SU(3) flavor symmetry for pions+kions
   to relate relevant decay modes.

Sometime, uncertainties could be cancelled

Guohuai Zhu, aXiv:1002.4518[hep-ph]

However, SU(3)<sub>F</sub> broken effect  $f_K/f_{\pi}$  sizable

For SU(3) symmetry broken effects, eg, Y.L.Wu,Y.F.Zhou, PRD72(2005)034037

# Nonperturbative Alternatives

■ QCD sumrules: [P.Ball,R.Zwicky, PRD'05,

T.Huang, Z.Li, F.Zuo, EPJC'09

Light Front Quark Model

[W.Jaus,PRD'90, H.Y.Cheng et al,PRD'97]

Renewed with SCET[C.D.Lu.W.Wang, Z.T.Wei, PRD'07]

Recently Extended to baryons:

Zentao,H.W.Ke, X.Q.Li,PRD'09, Y.M.Wang, Y.L.Sheng, C.D.Lu, PRD'09

Lattice QCD [to be reviewed by C.Liu]

# Pursue New Physics

#### So far, the best we can do:

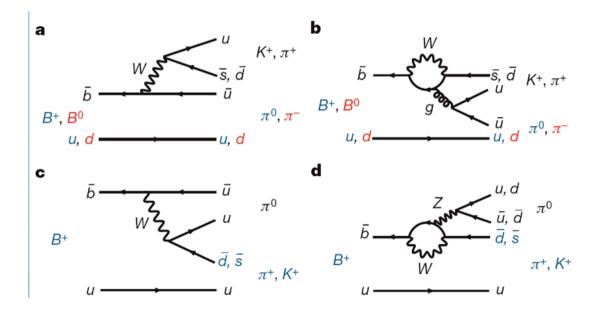
- Use as much form factor information from semileptonic decays as possible
- Use global fits which combine Factoriztion and SU(3) to look for interesting channels with large deviations from the SM
- Include THEORY uncertainties when discussing any deviations(model parameters and so on)
- Diagnose a new physics model if all the correlated deviations explained 
   constrained flavor structure 
   if it is accessible at LHC

### Examples: The $\pi$ K -CP puzzle

#### In the SM, the direct CP asymmetries in charge and neutral B decays

$$\begin{split} A_{\rm CP}(B^- \to K^- \pi^0) &\equiv \frac{\Gamma(B^- \to K^- \pi^0) - \Gamma(B^+ \to K^+ \pi^0)}{\Gamma(B^- \to K^- \pi^0) + \Gamma(B^+ \to K^+ \pi^0)} \\ A_{\rm CP}(\bar{B}^0 \to K^- \pi^+) &\equiv \frac{\Gamma(\bar{B}^0 \to K^- \pi^+) - \Gamma(B^0 \to K^+ \pi^-)}{\Gamma(\bar{B}^0 \to K^- \pi^+) + \Gamma(B^0 \to K^+ \pi^-)} \end{split}$$

Should be very close



### However, the measurements:

#### which established the difference:

$$\Delta A \equiv A_{\rm CP}(B^- \to K^- \pi^0) - A_{\rm CP}(\bar{B}^0 \to K^- \pi^+) = 0.164 \pm 0.037 \quad \text{at 4}$$

The average of BABAR, Belle, CDF & CLEO

4/17/10

5σ

Possible Implications

The mismatch may be due to:

Our limited understanding of QCD so far, say, strong phase.

Equally, new physics.

M. Peskin, Nature 452(2008)334

Model independent approach

The possible new physics effects termed

$$\begin{aligned} \mathcal{H}_{\text{eff}}^{\text{NP}} &= \frac{G_F}{\sqrt{2}} \sum_{a=u.d} |V_{tb} V_{ts}^*| e^{i\delta_S^q} \left[ C_{S1}^q O_{S1}^q + C_{S8}^q O_{S8}^q \right] + \text{h.c.} \,, \\ O_{S1}^u &= \bar{s}(1+\gamma_5) b \otimes \bar{u}(1+\gamma_5) u \,, \\ O_{S1}^d &= \bar{s}(1+\gamma_5) b \otimes \bar{d}(1+\gamma_5) d \,, \end{aligned} \qquad \begin{aligned} O_{S8}^u &= \bar{s}_i(1+\gamma_5) b_j \otimes \bar{u}_j(1+\gamma_5) u_i \,, \\ O_{S1}^d &= \bar{s}(1+\gamma_5) b \otimes \bar{d}(1+\gamma_5) d \,, \end{aligned}$$

#### i.e., scalar FCNCs

Scanning 12 correlated channels (CPA+Brs) all theoretical inputs and exp uncertainties included

# Application continued

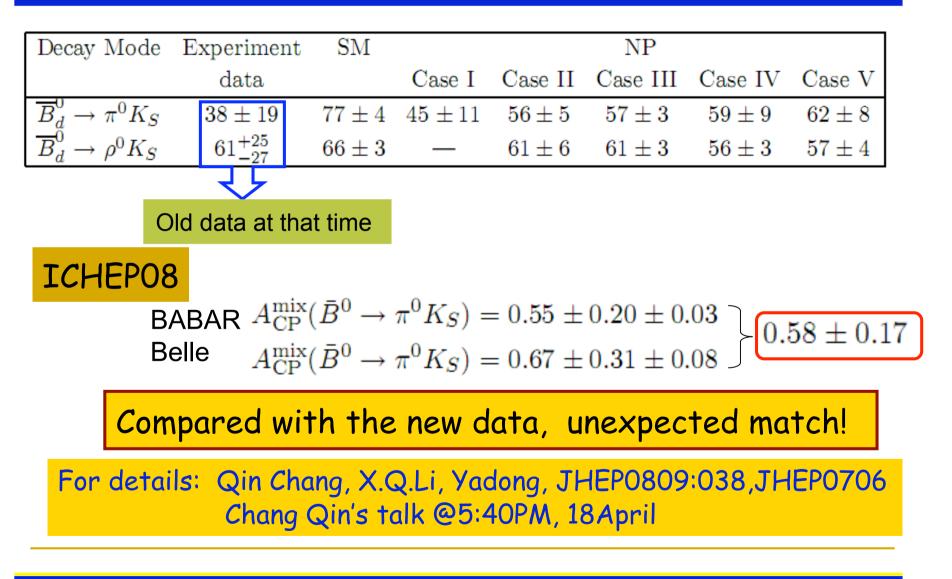
#### Constrained parameter space:

NP para.	Case I	Case II	Case III	Case IV	Case V
$C_{S1}^{u}(\times 10^{-3})$	$-41.6\pm13.4$	_	—	$25.8\pm8.4$	$-6.7\pm10.5$
$C_{S8}^u(\times 10^{-3})$	$38.7 \pm 18.2$	—	—	—	$16.0\pm7.1$
$\delta^u_S$	$99.5^\circ\pm6.1^\circ$	—	_	$107.0^\circ\pm11.5^\circ$	$73.0^\circ\pm23.8^\circ$
$C_{S1}^d(\times 10^{-3})$	_	$23.0\pm5.1$	$22.8\pm2.3$	$50.3 \pm 12.8$	$17.5\pm10.1$
$C_{S8}^d(\times 10^{-3})$	—	$-0.8\pm13.7$	—	—	$10.5\pm9.4$
$\delta^d_S$	_	$100.0^\circ\pm8.7^\circ$	$99.3^\circ\pm9.2^\circ$	$106.6^\circ\pm7.3^\circ$	$114.7^\circ\pm18.6^\circ$

Color singlet  $b \rightarrow sqq$  dominated + a nontrivial weak phase!

We leave the mixing induced CPA in  $B \to \pi^0 K_S$  and  $B \to \rho^0 K_S$  as our prediction

# Prediction(continued)



# Pursue in specified NP scenarios I

✓ Invisible Higgs in B→K vv-bar[Guohai et al, PRD'10]

✓ Family Non-universal Z' Models

To tackle A<sub>FB</sub>(q<sup>2</sup>) puzzle: C.W.Chiang, R.H.Li, C.D.Lu, arXiv:0911.2399 Q.Chang, X.Q. Li, Y.D.Yang JHEP in production

To tackle  $B_s - \overline{B}_s$ ,  $B_s \rightarrow \mu^+ \mu^{-}$ ,  $B_s \rightarrow X_s \mu^+ \mu^-$ Q.Chang, X.Q. Li, Y.D.Yang, JHEP1002(2010)

To tackle the abnormal large  $B \rightarrow \pi^0 \pi^0$ Q.Chang, X.Q. Li, Y.D.Yang,aXiv:1003.6051

P. Langacker, V.Barger... JHEP0912(2009)04, PRD80(09)055008

### Pursue in specified NP scenarios II

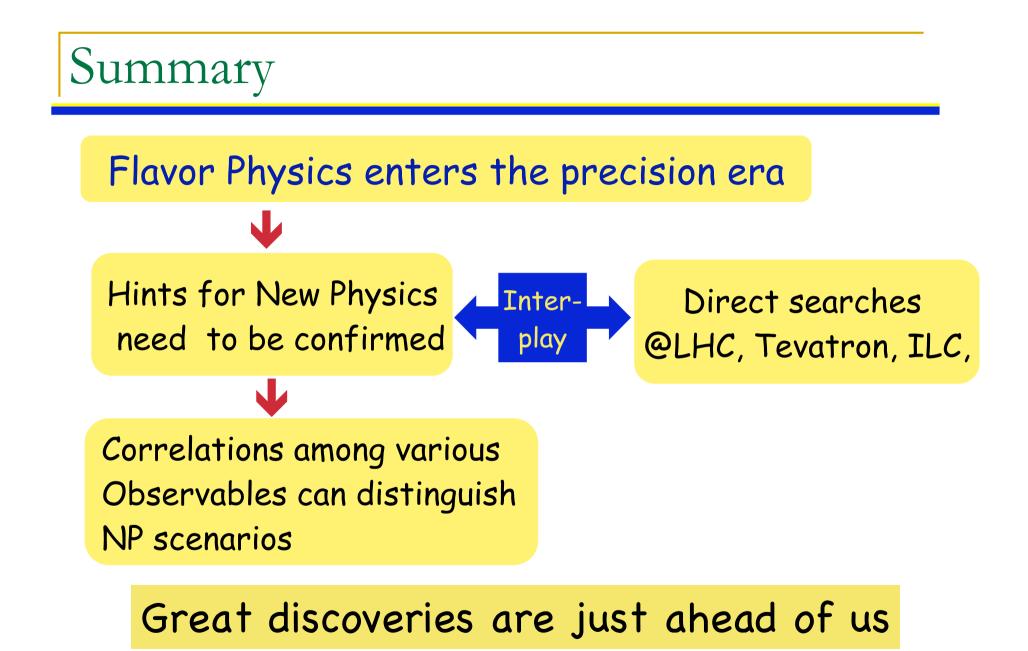
Most extensive studies from A.J. Buras grou

SUSY: AJB, Altmannshofer, S. Gori, P.Paradisi, D. Straub

LHT: AJB, M. Blanke, B. Duling, S. Recksiegel, C.Tarantino

RS: AJB Abrecht, K.Gemmler, A.Weiler

Munich gang



#### Which guy is the promised? Who will fly?

