

桂林, 2006 年 10 月 31 日

Determination of chiral couplings at NLO in $1/N_c$:

$$L_8^r(\mu)$$

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[ArXiv: hep-ph/0610290]

[ArXiv: hep-ph/0610304]

Introduction and motivation

Recurrent problem in non-perturbative QCD:

→ NEED FOR A QFT FOR MESONS (ρ , K^* , a_1 , f_0, \dots D, B)

A way to afford the problem:

- Large- N_c framework:
quark loops suppressed \Leftrightarrow meson loops suppressed
- Perturbative calculation within a $1/N_c$ expansion:
Systematic construction of a QFT for mesons
- QCD long-distance \Leftrightarrow Implementation of chiral symmetry:
We need a chiral theory for resonances ($R\chi T$)
- Interconnection of low/high-energy QCD through $R\chi T$:
Predictions for the $R\chi T$ and χPT coupling

Study of 2-point Green-functions

$$\Pi(t) = i \int d\mathbf{x}^4 e^{i\mathbf{q}\mathbf{x}} \left\langle T \left\{ \bar{q}\Gamma q(\mathbf{x}) \bar{q}\Gamma q(0)^\dagger \right\} \right\rangle, \text{ with } t=q^2$$

We focus the attention on the **SS-PP** with $I=1$ [ArXiv: hep-ph/0610290, hep-ph/0610304]


- At low energies: χ PT coupling L_8 (quark mass \leftrightarrow pGoldstone mass)


$$\Pi(t) = B_0^2 \left[\frac{2 F_\pi^2}{t} + 32 L_8^r(\mu) + \frac{\Gamma_8}{\pi^2} \left(1 - \ln \frac{-t}{\mu^2} \right) + \mathcal{O}(t) \right]$$


PROGRAM:

1. Construction of a chiral lagrangian for resonances ($R\chi T$)

2. Calculation of **form-factors** at LO 

3. Derivation of $\Pi(t)$ at NLO (dispersive relations) 

4. Matching OPE at NLO 

5. Recovery of χPT at NLO 

Reminder about
the Leading Order
in $1/N_c$

$$\Pi(t) = 2B_0^2 \left[\frac{F_\pi^2}{t} + \frac{8c_m^2}{M_S^2 - t} - \frac{8d_m^2}{M_P^2 - t} \right]$$

$$\Pi(t) \Big|_{\substack{R_{\pi T} \\ N_C \rightarrow \infty}} \xrightarrow{t \rightarrow -\infty} \Pi(t)^{\text{OPE}} = \sum_l \frac{\langle \mathcal{O}_{(2l)} \rangle}{(-t)^l}$$

$$\left. \begin{aligned} \langle \mathcal{O}_{(2)}^{\text{LR}} \rangle = 0 &\longrightarrow F_\pi^2 - 8c_m^2 + 8d_m^2 = 0 \\ \langle \mathcal{O}_{(4)}^{\text{LR}} \rangle \simeq 0 &\longrightarrow -8c_m^2 M_S^2 + 8d_m^2 M_P^2 \simeq 0 \end{aligned} \right\} \begin{array}{l} \text{WSRs} \\ \text{[Weinberg'67]} \end{array}$$

$$d_m^2 = \frac{F_\pi^2}{8} \left(\frac{M_S^2}{M_P^2 - M_S^2} \right), \quad c_m^2 = \frac{F_\pi^2}{8} \left(\frac{M_P^2}{M_P^2 - M_S^2} \right)$$

$$\Pi(t) = B_0^2 \left[\frac{2F_\pi^2}{t} + 32 \left(\frac{F_\pi^2}{16} \left\{ \frac{1}{M_S^2} + \frac{1}{M_P^2} \right\} \right) + \mathcal{O}(t) \right]$$

$L_8^{N_C \rightarrow \infty}$ At what μ ?

Program
for calculations
at NLO in $1/N_c$

1.

Ingredients of a Resonance Chiral Theory (R χ T)

[Ecker et al.'89], [Knecht & de Rafael'98]

...

- Large $N_c \rightarrow U(n_f)$ multiplets
- Goldstones from $S\chi SB$ (π, K, η_8, η_0)
- **MHA**: First resonance multiplets (V, A, S, P)
- **Chiral symmetry** invariance
- Just $\mathcal{O}(p^2)$ operators

$$\mathcal{L}_{R\chi T} = \mathcal{L}_{\chi PT}^{(2)} + \sum_{R_1} \mathcal{L}_{R_1} + \sum_{R_1, R_2} \mathcal{L}_{R_1, R_2} + \sum_{R_1, R_2, R_3} \mathcal{L}_{R_1, R_2, R_3} + \dots$$

$$\mathcal{L}_{\chi\text{PT}}^{(2)} = \frac{F^2}{4} \langle \mathbf{u}^\mu \mathbf{u}_\mu + \chi_+ \rangle$$

$$\mathcal{L}_{\chi\text{PT}}^{(2)}$$

[Weinberg'79]

[Gasser & Leutwyler'84]

[Gasser & Leutwyler'85]

$$\mathcal{L}_V[\pi, \mathbf{V}] = \frac{F_V}{2\sqrt{2}} \langle \mathbf{V}_{\mu\nu} \mathbf{f}_+^{\mu\nu} \rangle + \frac{i\mathbf{G}_V}{2\sqrt{2}} \langle \mathbf{V}_{\mu\nu} [\mathbf{u}^\mu, \mathbf{u}^\nu] \rangle$$

$$\mathcal{L}_S[\pi, \mathbf{S}], \quad \mathcal{L}_A[\pi, \mathbf{A}], \quad \mathcal{L}_P[\pi, \mathbf{P}]$$

$$\sum_{R_1} \mathcal{L}_{R_1}$$

[Ecker et al.'89]

...

couplings $\lambda_i^{RR}, \lambda_i^{RRR}$

$$\sum_{R_1, R_2} \mathcal{L}_{R_1, R_2} + \sum_{R_1, R_2, R_3} \mathcal{L}_{R_1, R_2, R_3}$$

... [Moussallam'95], [Knecht & Nyffeler'01]

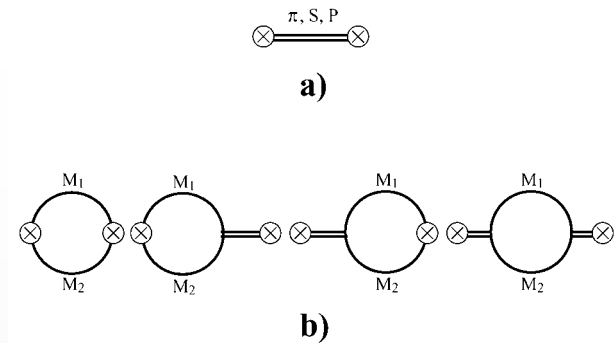
[Cirigliano et al.'06]

[Pich, Rosell & SC, forthcoming 06]

2.

Form-factors at LO in $1/N_c$:

Spectral function at NLO



$$\text{Im } \Pi(\mathbf{t})|_{\pi\pi} \sim |\mathcal{F}_{\pi\pi}(\mathbf{t})|^2 \xrightarrow{t \rightarrow \infty} 0 \quad \text{Vanishing Pion form factor}$$

[Brodsky & Lepage '79]



This same behaviour is demanded to $\text{Im } \Pi(\mathbf{t})|_{M_1, M_2}$

$$\text{Im } \Pi(\mathbf{t})|_{M_1, M_2} \propto |\mathcal{F}_{M_1, M_2}(\mathbf{t})|^2 \xrightarrow{t \rightarrow \infty} 0$$



Stringent constraints on the **form-factors**

and resonance couplings λ_j

in terms of M_S, M_P, M_V, M_A

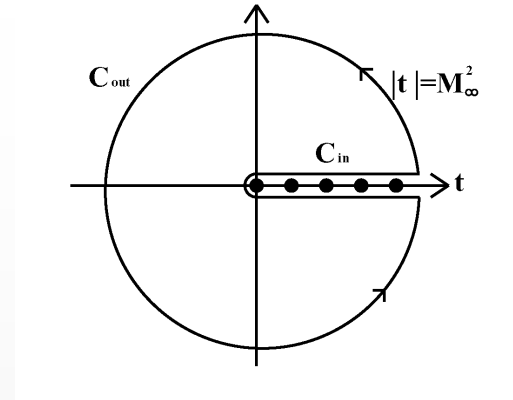
$$\left\{ \begin{array}{l} \mathcal{F}_{\pi\pi}^s(\mathbf{t}) = \frac{M_S^2}{M_S^2 - \mathbf{t}} \\ \mathcal{F}_{V\pi}^p(\mathbf{t}) = -2 \sqrt{1 - \frac{M_V^2}{M_A^2}} \frac{M_V^2 M_P^2}{(M_P^2 - \mathbf{t}) \mathbf{t}} \\ \dots \end{array} \right.$$

3.

Reconstruction of $\Pi(t)$ from $\text{Im } \Pi(t)$

through an unsubtracted dispersion relation:

$$\Pi(t) = \frac{1}{2\pi i} \oint \frac{dt'}{t'-t} \Pi(t') = \int_{\mathcal{R}^+} \frac{dt'}{t'-t} \frac{1}{\pi} \text{Im} \Pi(t')$$



At LO in $1/N_c$

$$\Pi(t) = 2B_0^2 \left[\frac{F_\pi^2}{t} + \frac{8c_m^2}{M_S^2 - t} - \frac{8d_m^2}{M_P^2 - t} \right]$$

[Rosell, Pich & SC'04], [Rosell *et al.*'05]

(UV) Renormalized counter-terms

One-loop diagrams

At NLO in $1/N_c$

$$\Pi(t) = 2B_0^2 \left[\frac{F^2}{t} + \frac{8c_m^{r^2}}{M_S^{r^2} - q^2} - \frac{8d_m^{r^2}}{M_P^{r^2} - q^2} \right] + \sum_{M_1, M_2} \Delta \Pi(t) \Big|_{M_1, M_2}$$

\uparrow
 $|\mathcal{F}_{M_1, M_2}(t)|^2$

4.

Matching $R\chi T$ and OPE up to NLO in $1/N_c$:

- High energy limit of one-loop contributions:

$$\sum_{M_1, M_2} \Delta\Pi(t)|_{M_1, M_2} \xrightarrow{t \rightarrow \infty} \frac{F^2}{t} (\delta_1) + \frac{F^2 M_S^2}{t^2} \left(\delta_2 + \tilde{\delta}_2 \ln \frac{-t}{M_S^2} \right) + \mathcal{O}\left(\frac{1}{t^3}\right)$$

$f(M_S, M_P, M_V, M_A)$

- Matching the whole $\Pi(t)$ to OPE leads to

-No $\frac{1}{t^2} \ln \frac{-t}{M_S^2}$ term $\rightarrow \tilde{\delta}_2 = 0$

-1st and 2nd WSRs modified by terms NLO in $1/N_c$:

$$\left. \begin{aligned} \langle \mathcal{O}_{(2)}^{LR} \rangle = 0 &\longrightarrow F_\pi^2 (1 + \delta_1) - 8 c_m^{r^2} + 8 d_m^{r^2} = 0 \\ \langle \mathcal{O}_{(4)}^{LR} \rangle \simeq 0 &\longrightarrow -8 c_m^{r^2} M_S^{r^2} + 8 d_m^{r^2} M_P^{r^2} + F_\pi^2 M_S^2 \delta_2 = 0 \end{aligned} \right\}$$

$$\longrightarrow d_m^{r^2} = \frac{F^2}{8} \left(\frac{M_S^{r^2}}{M_P^{r^2} - M_S^{r^2}} \right) \times [1 + \delta_1 - \delta_2], \dots$$

5.

Recovery of χ PT at the one-loop level:

- Chiral symmetry ensures the right low-energy dynamics

(governed by χ PT):

- Proper non-analytic $\ln(-t)$ structure
- Proper running of the LECs

$$\Pi(t)^{\text{R}\chi\text{T}} = \mathbf{B}_0^2 \left[\frac{2 F_\pi^2}{t} + 32 \bar{\mathbf{L}}_8 + \frac{\Gamma_8}{\pi^2} \left(1 - \ln \frac{-t}{M_S^2} \right) + \mathcal{O}(t) \right]$$



$f(M_S^r, M_P^r, M_V)$

Prediction for $L_8^r(\mu)$ at any μ

E.g., for $\mu_0=770$ MeV: $10^3 \cdot L_8^r(\mu_0) = 0.6 \pm 0.4$

Conclusions

- Chiral symmetric resonance lagrangian ($R\chi T$):

Essential to recover low-energy QCD (χPT)

- $1/N_c$ expansion:

Systematic expansion to compute loops with heavy particle

- Matching to high-energy QCD:

Predictions for (renormalized) $R\chi T$ and χPT

- WE WANT A REAL QFT FOR MESONS

(not just narrow-width or Breit-Wigner ansate) !!!