# Galaxy Rotation Curves from String Theory 

Yeuk－Kwan Edna Cheung 张若笉<br>Dept．of Physics，Nanjing U．<br>CPS2006，Autumn Meeting

- Work in Progress with:
- Hsien-Chung Kao, Taiwan Normal University
- Konstantin Savvidy, Nanjing University
- Feng Xu
- String Theory is:
- A theory of Everything?
too many assumptions; everything seems possible!
- A theory of Quantum Gravity?
provides the needed fundamental theory for cosmology? yet to make contact with observation.
- A theory good for Nothing? spectacular success in mathematics over the past decade. no experimental support whatsoever...


## Pauli would have said, "It is not even WRONG!"

Breaking the "symmetry" between right and wrong...

Getting in touch with data is very rewarding by itself!

## The Missing Mass Problem:

Young and bright stars in a spiral galaxy lie on a thin stellar disk. They execute circular orbits around the center of the galaxy.


## What Nature does:



## Dark Energy:



DM \& DE are the two main roadblocks on our path to a comprehensive fundamental theory of Nature.

- Dark Matter:
- baryonic or non-baryonic;
- well-founded or exotic;
- MOND: modification of Newtonian dynamics at large scale;
- existent fields/long range force from String Theory; exploit their low energy implications.
point particle and gauge field:

$$
\int A \cdot d X \quad F=d A
$$

string and its gauge potential:

$$
\int B \cdot d S \quad H=d B
$$

particle worldline
If strings are indeed fundamental objects, then $B_{\mu \nu}$ will play a role as fundamental as $A_{\mu}$ does.

- plane-polarized gravitational fields:

centre of mass of the closed string follows the geodesic:

$$
\begin{aligned}
u & =u_{0}+H p^{+} \tau \\
a & =-\lambda+\rho e^{+i H p^{+} \tau} \\
\bar{a} & =-\bar{\lambda}+\bar{\rho} e^{-i H p^{+} \tau}
\end{aligned}
$$

A short string in
a Landau orbit with
Larmore frequency $=\mathrm{Hp}^{+}$

In a leap of faith...

- If matter is indeed made of strings, they will all be charged under this "gravi-magnetic" field.
- In the presence of such background field, galaxies will execute Landau orbits.
- provides the extra centripetal force, which would otherwise be attributed to extra mass:

$$
m \frac{v^{2}}{r}=Q H_{z} v+\frac{G_{N} M m}{r^{2}}
$$

extra mass $\Rightarrow$ Dark Matter

## Parametric Modeling of the Mass Distribution

- Van Der Kruit \& Searle’s Formula:

$$
\rho(r, z)=\rho_{0} \exp \left(-\frac{r}{R_{d}}\right) \operatorname{sech}^{2}\left(\frac{6 z}{R_{d}}\right)
$$

for the visible stellar disk and spheroid.

- introduce three parameters: $\Omega, \rho$, and $R_{d}$.

$$
v^{2}=R_{o b s} \Omega v+R_{o b s} \rho \tilde{E}(\tilde{r})
$$

where $\tilde{E}(\tilde{r})$ is dimensionless.

- data: http://www.astro.umontreal.ca/fantomm/sings/index.htm










$$
\rho(r, z)=\rho_{0} \exp \left(-\frac{r}{R_{d}}\right) \operatorname{sech}^{2}\left(\frac{6 z}{R_{d}}\right)
$$

| Galaxy | Rd | rho | Mass (M sun)2 |
| :---: | :---: | :---: | :---: |
| ngc0628 | 29.67 | 847.01 | 4.19E+11 |
| ngc0925 | 285.68 | 128.63 | $5.90 \mathrm{E}+12$ |
| ngc2403 | 31.19 | 142.06 | 7.77E+10 |
| ngc3031 | 8.47 | 41.92 | $1.69 \mathrm{E}+9$ |
| ngc3184 | 69.54 | 327.88 | 8.91E+11 |
| ngc3198 | 60.41 | 408.20 | 8.37E+11 |
| ngc3521 | 26.67 | 2197.00 | $8.78 \mathrm{E}+11$ |
| ngc4236 | 350.00 | 6.01 | 4.13E+11 |
| ngc4321 | 27.37 | 1360.15 | $5.72 \mathrm{E}+11$ |
| ngc4536 | 45.90 | 551.03 | $6.52 \mathrm{E}+11$ |
| ngc4569 | 16.78 | 789.54 | $1.24 \mathrm{E}+11$ |
| ngc4579 | 43.88 | 2171.41 | $2.35 \mathrm{E}+12$ |
| ngc5055 | 30.59 | 1914.81 | $1.01 \mathrm{E}+12$ |
| ngc5194 | 24.35 | 507.39 | $1.69 \mathrm{E}+11$ |
| ngc5713 | 20.00 | 454.78 | $1.02 \mathrm{E}+11$ |
| ngc6946 | 56.81 | 293.73 | $5.32 \mathrm{E}+1 \mathrm{l}$ |

## Frank C. van den Bosch et al, astro-ph/99| I 372

Table 2. Parameters of best fits to HI surface brightness.

| Galaxy (1) | $\mathrm{M}_{\odot} \mathrm{pc}^{-2}$ <br> (2) | $h_{70}^{-1} \mathrm{kpc}$ (3) | $\beta$ (4) | $R_{c}$ $h_{70}^{-1} \mathrm{kpc}$ (5) | $\begin{gathered} \log \left(M_{\mathrm{HI}}\right) \\ h_{70}^{-2} \mathrm{M}_{\odot} \end{gathered}$ <br> (6) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| F563-1 | 8.59 | 10.63 | 0.20 | 26.37 | 9.644 |
| F568-1 | 4.55 | 1.97 | 3.43 | 16.98 | 9.674 |
| F568-3 | 11.52 | 3.46 | 1.78 | 19.45 | 9.524 |
| F568-V1 | 11.55 | 5.24 | 1.39 | 15.91 | 9.464 |
| F574-1 | 2.16 | 3.13 | 3.51 | 18.71 | 9.649 |
| F583-1 | 9.38 | 2.77 | 2.09 | 16.18 | 9.401 |
| NGC 247 | 4.24 | 0.56 | 7.89 | 8.63 | 8.912 |
| DDO 154 | 14.38 | 1.53 | 0.52 | 6.17 | 8.383 |
| NGC 3109 | 8.28 | 3.08 | 0.32 | 12.92 | 8.713 |

Note. - Column (1) lists the name of the galaxy. Columns (2) through (5) list the best fitting parameters for the HI surface density, and column (6) lists the corresponding HI mass.

Table 3. Parameters of fits to rotation curves.

| Galaxy <br> $(1)$ | Model | $\alpha$ | $c$ | $V_{200}$ | $\Upsilon_{B}$ | $f_{\text {bar }}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $(2)$ | $(3)$ | $(4)$ | $(5)$ | $(6)$ | $(7)$ |
|  |  |  |  |  |  |  |
| F563-1 | BF | 2.00 | 5.2 | 73.5 | 0.0 | 0.039 |
|  | BF | 1.97 | 5.8 | 64.0 | 6.2 | 0.369 |
| F568-3 | BF | 1.18 | 3.4 | 127.7 | 0.5 | 0.010 |
| F568-V1 | BF | 0.47 | 15.6 | 91.6 | 0.9 | 0.023 |
| F574-1 | BF | 0.26 | 8.6 | 118.3 | 1.0 | 0.018 |
|  | a | 1.30 | 8.6 | 76.4 | 1.0 | 0.067 |
|  | b | 0.26 | 8.6 | 55.7 | 6.0 | 0.537 |
| F583-1 | c | 0.80 | 2.0 | 278.8 | 1.0 | 0.001 |
| NGC 247 | BF | 0.00 | 20.6 | 65.7 | 0.0 | 0.035 |
| DDO 154 | BF | 1.02 | 7.2 | 93.1 | 1.0 | 0.011 |
| NGC 3109 | BF | 0.00 | 14.7 | 44.0 | 0.0 | 0.011 |
|  | BF | 0.00 | 10.2 | 101.6 | 0.0 | 0.002 |

Note. - Column (1) lists the name of the galaxy. Columns (2) lists the ID of the model, with 'BF' indicating the best-fit model (i.e., the one that minimizes $\chi_{\text {vel }}^{2}$ ). For F574-1 three additional models are listed ( $\mathrm{a}, \mathrm{b}$, and c ) all of which fall within the 68.3 confidence level of the BF-model (see contour plots in Figure 4). Columns (3) through (5) list parameters of the model: $c, \Upsilon_{B}$ (in $h_{70} \mathrm{M}_{\odot} / \mathrm{L}_{\odot}$ ), and $V_{200}$ (in $\mathrm{km} \mathrm{s}^{-1}$ ). Finally, column (7) gives the resulting baryon fraction $f_{\text {bar }}=\left(M_{\text {gas }}+M_{\text {stars }}\right) / M_{200}$

## Comments:

- only 3 parameters vs the usual 8
- masses of the galaxies obtained
- cross-check with photometric method In progress
- effective field theory...
In progress

I have pushed the limits of the model...

## Perhaps a happy ending:

- Theoretically well-motivated, (very ordinary) Dark Matters:
- black holes
- small stars (not burning hydrogen)
- A little bit of String Gauge Fields
in a happy union!


## fin

## Thank You!

